# SHARP SOLVABILITY CONDITIONS FOR A FOURTH ORDER EQUATION WITH PERTUBATION 

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Abstract. Let $B$ be the unit ball of $\mathbb{R}^{n}, n \geq 5$, and $\rho: \mathbb{R} \rightarrow \mathbb{R}$ a smooth function. We consider the following critical problem

$$
\begin{cases}\Delta^{2} u=|u|^{\frac{8}{n-4}} u+\rho(u) & \text { in } B \\ u \not \equiv 0 & \\ u=\frac{\partial u}{\partial n}=0 & \text { on } \partial B\end{cases}
$$

We give sufficient conditions for the existence of solutions to this problem. These conditions are close to be sharp, as we prove by considering the problem on arbitrary small balls.

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## 1. Introduction and statement of the results

Let $n \geq 5$. We denote by $B(0, r) \subset \mathbb{R}^{n}$ the $n$-dimensional ball of radius $r>0$ and centered at 0 . Let $\rho \in C^{\infty}(\mathbb{R})$ be a smooth function. For $r>0$, we are interested in finding solutions $u \in C^{4}(\bar{B}(0, r))$ to the following problem:

$$
\begin{cases}\Delta^{2} u=|u|^{2^{\sharp}-2} u+\rho(u) & \text { in } \bar{B}(0, r)  \tag{r}\\ u \neq 0 & \text { on } \partial \bar{B}(0, r) .\end{cases}
$$

where $\Delta=-\sum \frac{\partial^{2}}{\partial x_{i}^{2}}$ is the Laplacian with the minus sign convention, $\frac{\partial}{\partial n}$ denotes the normal derivative with respect to the unit outward vector $\vec{n}$, and $2^{\sharp}=\frac{2 n}{n-4}$ is critical from the viewpoint of Sobolev embeddings. More precisely, for $\Omega \subset \mathbb{R}^{n}$ an open subset, we denote by $H_{2,0}^{2}(\Omega)$ the standard Sobolev space of second order, that is the completion of $C_{c}^{\infty}(\Omega)$, the set of smooth compactly supported functions in $\Omega$, with respect to the norm

$$
\|u\|_{H_{2,0}^{2}(\Omega)}=\sqrt{\int_{\Omega}(\Delta u)^{2} d x} .
$$

It follows from the Sobolev embedding theorem that $H_{2,0}^{2}(\Omega)$ is continuously embedded in $L^{q}(\Omega)$ for $1 \leq q \leq 2^{\sharp}$, and that this embedding is compact if and only if $1 \leq q<2^{\sharp}$. This lack of compactness is one of the main difficulties attached to problem $\left(E_{r}\right)$. Moreover, see $[\mathrm{Osw}]$, it can be shown that $\left(E_{r}\right)$ has no positive solution if $\rho \equiv 0$.

[^0]This type of problem was first studied by Brézis and Nirenberg. In [BrNi], Brézis and Nirenberg studied the existence of solutions to the elliptic problem

$$
\begin{cases}\Delta u=u^{\frac{n+2}{n-2}}+\rho(u) & \text { in } B_{1} \\ u>0 & \text { in } B_{1} \\ u=0 & \text { on } \partial B_{1}\end{cases}
$$

Using test-functions arguments and the moutain-pass lemma of Ambrosetti and Rabinowitz [AmRa], they prove that $\left(E^{\prime}\right)$ possesses a solution if $\rho(0)=\rho^{\prime}(0)=0$ and $\int_{0}^{+\infty} \rho(s) s^{-\frac{n}{n-2}} d s>0$. Later on, in view of a nonexistence result of Adimurthi and Yadava in the absence of the above condition, Brézis raised the following question: is the preceding condition a necessary and sufficient condition for the existence of positive solutions for $\left(E^{\prime}\right)$ when $\rho$ is compactly supported? A first step in answering this question was carried out by Adimurthi and Yadava [AdYa]. When $n \geq 7$, they prove that there is a specific class of functions $\rho$ for which ( $E^{\prime}$ ) has a solution if and only if $\int_{0}^{+\infty} \rho(s) s^{-\frac{n}{n-2}} d s \geq 0$. Adimurthi, Mancini and Sandeep [AMS] came back to this problem for a fairly general class of functions $\rho$. They introduced a new set of conditions for the solvability of $\left(E^{\prime}\right)$ in higher dimensions. In particular, using blow-up analysis they showed that there exist functions $\rho$ such that $\int_{0}^{+\infty} \rho(s) s^{-\frac{n}{n-2}} d s=0$, and problem $\left(E^{\prime}\right)$ does not have solutions on arbitrary small balls. We refer to [AMS] for more details.

Let us now return to the study of $\left(E_{r}\right)$. There has been considerable interest in higher order operators since the pioneering work of Chang, Gursky and Yang concerning the Paneitz operator on Riemannian manifolds. We refer for instance to [Cha] for a general survey on such operators. We refer also to [EFJ], [PuSe], [VdV] in the Euclidean context, and [DHL], [HeRo] in the Riemannian context.

In this paper we address questions similar to the ones addressed in [AMS], but concerning the bi-harmonic operator. To be more precise we define

$$
\begin{align*}
& I_{1}(\rho)=\int_{0}^{+\infty} \rho(s) s^{-\frac{n}{n-4}} d s, \quad I_{2}(\rho)=\int_{0}^{+\infty} \rho(s) s^{-\frac{n-2}{n-4}} d s \\
& I_{3}(\rho)=\int_{0}^{+\infty} r^{-\frac{2 n-4}{n-4}}\left[\int_{0}^{r} t^{\frac{2}{n-4}}\left(\int_{t}^{+\infty} \rho(s) s^{-\frac{2 n-4}{n-4}} d s\right) d t\right]^{2} d r  \tag{1}\\
& +\frac{(n-4)^{4}}{4 n(n+2)} \int_{0}^{+\infty} \rho(t) \frac{1}{t} d t
\end{align*}
$$

when these quantities make sense. We say that $u \in C^{4}(\bar{B}(0, r))$ is a solution of small energy for $\left(E_{r}\right)$ if it is a solution of $\left(E_{r}\right)$ satisfying

$$
\frac{1}{2} \int_{B(0, r)}(\Delta u)^{2} d x-\frac{1}{2^{\sharp}} \int_{B(0, r)}|u|^{2^{\sharp}} d x-\int_{B(0, r)} \tilde{\rho}(u) d x<\frac{2}{n K_{0}^{\frac{n}{4}}},
$$

where $\tilde{\rho}(r)=\int_{0}^{r} \rho(t) d t$ for $r \in \mathbb{R}$, and $K_{0}>0$ is the best constant in the second order Sobolev inequality. Namely

$$
\frac{1}{K_{0}}=\inf \frac{\int_{\mathbb{R}^{n}}(\Delta u)^{2} d x}{\left(\int_{\mathbb{R}^{n}}|u|^{2^{\sharp}} d x\right)^{\frac{2}{2 \sharp}}},
$$

where the infimum is taken over the nonzero compactly supported functions in $\mathbb{R}^{n}$. We assume in what follows that

$$
\begin{align*}
& \rho(0)=\rho^{\prime}(0)=0, \rho^{\prime} \text { is bounded } \\
& \exists b>\frac{2}{n-4} \text { such that }|\rho(s)| \leq C|s|^{-b} \text { for all } s \neq 0
\end{align*}
$$

Our main result is the following:
Theorem 1.1. Assume that $n \geq 13$ and that $\left(H_{\rho}\right)$ holds. If $I_{1}(\rho)>0$, or if $I_{1}(\rho)=0$ and $I_{2}(\rho)<0$, or if $I_{1}(\rho)=I_{2}(\rho)=0$ and $I_{3}(\rho)>0$, then $\left(E_{r}\right)$ has a radially symmetrical solution of small energy for all $r>0$. Conversely, if $\left(E_{r}\right)$ has a radially symmetrical solution of small energy for all $r>0$, then $I_{1}(\rho) \geq 0$ with the additional properties that if $I_{1}(\rho)=0$, then $I_{2}(\rho) \leq 0$, and if $I_{1}(\rho)=I_{2}(\rho)=0$, then $I_{3}(\rho) \geq 0$.

When we deal with an arbitrary subset of $\mathbb{R}^{n}$, the existence part still holds, but the solutions are not necessarily radially symmetrical. The paper is divided as follows. Section 2 is devoted to test-functions estimates. We prove the existence part of theorem 1.1 in section 3 . Sections 4, 5 are devoted to the blow-up analysis attached to our problem, and to the proof the second part of theorem 1.1. In section 6 , we prove a spectral result we need in section 5 . Extensions of theorem 1.1 to the case of a smooth open subset of $\mathbb{R}^{n}$ and to smaller dimensions are discussed at the end of sections 3 and 5 .

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## 2. Test-functions estimates

We consider a function $\rho \in C^{\infty}(\mathbb{R})$ satisfying the following conditions:

$$
\begin{align*}
& \rho(0)=\rho^{\prime}(0)=0, \rho^{\prime} \text { is bounded } \\
& \exists b>\frac{2}{n-4} \text { such that }|\rho(s)| \leq C|s|^{-b} \text { for all } s \neq 0 \tag{2}
\end{align*}
$$

We also define $\tilde{\rho}(r)=\int_{0}^{r} \rho(t) d t, r>0$. We denote by $B$ the unit ball of $\mathbb{R}^{n}$, and for $\alpha>0$, we consider the following functional

$$
J_{\alpha}(u)=\frac{1}{2} \int_{B}(\Delta u)^{2} d x-\frac{1}{2^{\sharp}} \int_{B}|u|^{2^{\sharp}} d x-\alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{u}{\alpha}\right) d x,
$$

where $u \in H_{2,0}^{2}(B)$. We define the function $U \in H_{2,0}^{2}\left(\mathbb{R}^{n}\right)$ by

$$
\begin{equation*}
U(x)=\left(\frac{a_{n}^{2}}{a_{n}^{2}+|x|^{2}}\right)^{\frac{n-4}{2}} \tag{3}
\end{equation*}
$$

where $x \in \mathbb{R}^{n}$, and $a_{n}=\sqrt[4]{n(n-4)\left(n^{2}-4\right)}$. It is easily checked that $U$ verifies $\Delta^{2} U=U^{2^{\sharp}-1}$. Moreover, $U$ is an extremal for the second order Sobolev inequality

$$
\begin{equation*}
\frac{1}{K_{0}}=\inf _{u \in H_{2,0}^{2}\left(\mathbb{R}^{n}\right)-\{0\}} \frac{\int_{\mathbb{R}^{n}}(\Delta u)^{2} d x}{\left(\int_{\mathbb{R}^{n}}|u|^{2^{\sharp}} d x\right)^{\frac{2}{2 \sharp}}}, \tag{4}
\end{equation*}
$$

The value of $K_{0}>0$ and the extremals for (4) are explicitely known. They have been computed by Lieb [Lie], Lions [Lio], and Edmunds-Fortunato-Janelli [EFJ]. For any $\varepsilon>0$, we define

$$
\begin{equation*}
U_{\varepsilon}(x)=\varepsilon^{-\frac{n-4}{2}} U\left(\frac{x}{\varepsilon}\right)=\left(\frac{a_{n}^{2} \varepsilon}{a_{n}^{2} \varepsilon^{2}+|x|^{2}}\right)^{\frac{n-4}{2}} \tag{5}
\end{equation*}
$$

where $x \in \mathbb{R}^{n}$, and $a_{n}$ is as above. Then

$$
\begin{equation*}
\Delta^{2} U_{\varepsilon}=U_{\varepsilon}^{2^{\sharp}-1} . \tag{6}
\end{equation*}
$$

From now on, if $R>0$ and if $h: B(0, R) \rightarrow \mathbb{R}$ is a radially symmetrical function, we write $h(r)=h(|x|)$, where $x \in B(0, R)$ and $|x|=r$. For $\alpha, \varepsilon>0$, we consider the unique radially symmetrical function $v_{\varepsilon, \alpha} \in C^{4}(\bar{B})$ solution of the problem:

$$
\begin{cases}\Delta^{2} v_{\varepsilon, \alpha}=\alpha^{2^{\sharp}-1} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) & \text { in } B  \tag{7}\\ v_{\varepsilon, \alpha}+U_{\varepsilon}=\frac{\partial\left(v_{\varepsilon, \alpha}+U_{\varepsilon}\right)}{\partial n}=0 & \text { on } \partial B .\end{cases}
$$

This function is explicitely known. We have that

$$
\begin{align*}
& v_{\varepsilon, \alpha}(r)=-U_{\varepsilon}(1)-\frac{C_{\varepsilon, \alpha}}{2 n}\left(1-r^{2}\right)  \tag{8}\\
& -\alpha^{2^{\sharp}-1} \int_{r}^{1} t^{1-n}\left[\int_{0}^{t} s^{n-1}\left[\int_{0}^{s} u^{1-n}\left\{\int_{0}^{u} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) v^{n-1} d v\right\} d u\right] d s\right] d t
\end{align*}
$$

where

$$
C_{\varepsilon, \alpha}=-n \frac{\partial U_{\varepsilon}}{\partial n}(1)-n \alpha^{2^{\sharp}-1} \int_{0}^{1} s^{n-1}\left[\int_{0}^{s} u^{1-n}\left\{\int_{0}^{u} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) v^{n-1} d v\right\} d u\right] d s
$$

In the sequel, $a_{\varepsilon, \alpha}=O\left(b_{\varepsilon, \alpha}\right)$ means that there exists $C>0$ independant of $\varepsilon>0$ and $\alpha \in(0,1]$ such that $\left|a_{\varepsilon, \alpha}\right| \leq C\left|b_{\varepsilon, \alpha}\right|$. We write $a_{\varepsilon, \alpha}=o\left(b_{\varepsilon, \alpha}\right)$ if for any $\eta>0$, there exists $\varepsilon_{0}>0$ such that $\left|a_{\varepsilon, \alpha}\right| \leq \eta\left|b_{\varepsilon, \alpha}\right|$ for all $\varepsilon \in\left(0, \varepsilon_{0}\right)$ and all $\alpha \in(0,1]$. With (2), it follows from (7) that

$$
\begin{equation*}
\left\|v_{\varepsilon, \alpha}\right\|_{H_{2}^{2}(B)}=o(1) . \tag{9}
\end{equation*}
$$

Here and in what follows, $H_{k}^{p}(\Omega)$ denotes the Sobolev space of functions $u \in L^{p}(\Omega)$ such that $\nabla^{i} u \in L^{p}(\Omega)$ for $i=1 \ldots k$, where $\Omega$ is an open subset of $\mathbb{R}^{n}$.

This section is devoted to finding estimate on

$$
J_{\alpha}\left(u_{\varepsilon, \alpha}\right)=\frac{1}{2} \int_{B}\left(\Delta u_{\varepsilon, \alpha}\right)^{2} d x-\frac{1}{2^{\sharp}} \int_{B}\left|u_{\varepsilon, \alpha}\right|^{2^{\sharp}} d x-\alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{u_{\varepsilon, \alpha}}{\alpha}\right) d x,
$$

where $u_{\varepsilon, \alpha}=U_{\varepsilon}+v_{\varepsilon, \alpha}+W_{\varepsilon, \alpha}$, and $W_{\varepsilon, \alpha} \in H_{2,0}^{2}(B)$ is assumed to be such that

$$
\begin{equation*}
\left\|W_{\varepsilon, \alpha}\right\|_{H_{2,0}^{2}(B)}=o\left(\varepsilon^{\frac{n-4}{2}}\right)+o\left(\left\|\Delta v_{\varepsilon, \alpha}\right\|_{2}\right) \tag{10}
\end{equation*}
$$

Here and in the sequel, $\|\cdot\|_{p}$ denotes the $L^{p}-$ norm for all $p \geq 1$.
Step 1: We first claim that

$$
\begin{equation*}
\int_{B} U_{\varepsilon}^{2^{\sharp}-1}\left|v_{\varepsilon, \alpha}\right| d x=o\left(\varepsilon^{\frac{n-4}{2}}\right), \int_{B} U_{\varepsilon}^{2^{\sharp}-2} v_{\varepsilon, \alpha}^{2} d x=o\left(\varepsilon^{n-4}+\left\|\Delta v_{\varepsilon, \alpha}\right\|_{2}^{2}\right) \tag{11}
\end{equation*}
$$

We prove the claim. We let $v_{\varepsilon, \alpha}^{+}, v_{\varepsilon, \alpha}^{-} \in H_{2,0}^{2}(B) \cap C^{4}(\bar{B})$ be radially symmetrical functions such that

$$
\left\{\begin{array}{ll}
\Delta^{2} v_{\varepsilon, \alpha}^{+}=\alpha^{2^{\sharp}-1} \rho^{+}\left(\frac{U_{\varepsilon}}{\alpha}\right) & \text { in } B \\
v_{\varepsilon, \alpha}^{+}=\frac{\partial v_{\varepsilon, \alpha}^{+}}{\partial n}=0 & \text { on } \partial B
\end{array},\left\{\begin{array}{ll}
\Delta^{2} v_{\varepsilon, \alpha}^{-}=\alpha^{2^{\sharp}-1} \rho^{-}\left(\frac{U_{\varepsilon}}{\alpha}\right) & \text { in } B \\
v_{\varepsilon, \alpha}^{+}=\frac{\partial v_{\varepsilon, \alpha}^{+}}{\partial n}=0 & \text { on } \partial B
\end{array},\right.\right.
$$

where $\rho^{+}(s)=\max \{\rho(s), 0\}$ and $\rho^{-}(s)=\max \{-\rho(s), 0\}$ for all $s \in \mathbb{R}$. As stated in Boggio [Bog] (see also Grunau-Sweers [GrSw]), the Green's function on the ball for the bi-harmonic operator with Dirichlet boundary condition is positive. It then follows that $v_{\varepsilon, \alpha}^{+}$and $v_{\varepsilon, \alpha}^{-}$are nonnegative. We define

$$
\begin{equation*}
T_{\varepsilon}(x)=\left(U_{\varepsilon}(1)-\frac{1}{2} \frac{\partial U_{\varepsilon}}{\partial n}(1)\right)+\frac{1}{2} \frac{\partial U_{\varepsilon}}{\partial n}(1)|x|^{2} \tag{12}
\end{equation*}
$$

for all $x \in B$. Clearly,

$$
\Delta^{2} T_{\varepsilon}=0 \text { in } B, \quad \text { and } T_{\varepsilon}=U_{\varepsilon}, \frac{\partial T_{\varepsilon}}{\partial n}=\frac{\partial U_{\varepsilon}}{\partial n} \text { on } \partial B
$$

Similarly, $T_{\varepsilon}>0$ and $U_{\varepsilon}-T_{\varepsilon} \in H_{2,0}^{2}(B)$. Now, integrating by parts, we get that

$$
\begin{aligned}
\int_{B} U_{\varepsilon}^{2^{\sharp}-1} v_{\varepsilon, \alpha}^{+} d x & =\int_{B} \Delta^{2} U_{\varepsilon} v_{\varepsilon, \alpha}^{+} d x=\int_{B} \Delta U_{\varepsilon} \Delta v_{\varepsilon, \alpha}^{+} d x \\
& =\int_{B} \Delta\left(U_{\varepsilon}-T_{\varepsilon}\right) \Delta v_{\varepsilon, \alpha}^{+} d x+\int_{B} \Delta T_{\varepsilon} \Delta v_{\varepsilon, \alpha}^{+} d x \\
& =\int_{B} \Delta\left(U_{\varepsilon}-T_{\varepsilon}\right) \Delta v_{\varepsilon, \alpha}^{+} d x=\int_{B}\left(U_{\varepsilon}-T_{\varepsilon}\right) \Delta^{2} v_{\varepsilon, \alpha}^{+} d x \\
& =\alpha^{2^{\sharp}-1} \int_{B} U_{\varepsilon} \rho^{+}\left(\frac{U_{\varepsilon}}{\alpha}\right) d x+O\left(\varepsilon^{\frac{n-4}{2}} \alpha^{2^{\sharp}-1} \int_{B} \rho^{+}\left(\frac{U_{\varepsilon}}{\alpha}\right) d x\right)
\end{aligned}
$$

With (2), it comes that for all $\nu \in(0,1)$, there exists $C_{\nu}>0$ such that $\rho^{+}(s) \leq C_{\nu} s^{\nu}$ for all $s>0$, and we get that

$$
\int_{B} U_{\varepsilon}^{2^{\sharp}-1} v_{\varepsilon, \alpha}^{+} d x=o\left(\varepsilon^{\frac{n-4}{2}}\right) .
$$

Similarly,

$$
\int_{B} U_{\varepsilon}^{2^{\sharp}-1} v_{\varepsilon, \alpha}^{-} d x=o\left(\varepsilon^{\frac{n-4}{2}}\right) .
$$

Now, using that $v_{\varepsilon, \alpha}=v_{\varepsilon, \alpha}^{+}-v_{\varepsilon, \alpha}^{-}-T_{\varepsilon}$, and that $v_{\varepsilon, \alpha}^{+}, v_{\varepsilon, \alpha}^{-}, T_{\varepsilon} \geq 0$, we get that

$$
\int_{B} U_{\varepsilon}^{2^{\sharp}-1}\left|v_{\varepsilon, \alpha}\right| d x=o\left(\varepsilon^{\frac{n-4}{2}}\right) .
$$

This proves the first equation in (11). Now, with Hölder's and Young's inequalities, we get

$$
\begin{aligned}
\int_{B} U_{\varepsilon}^{2^{\sharp}-2} v_{\varepsilon, \alpha}^{2} d x & \leq\left(\int_{B} U_{\varepsilon}^{2^{\sharp}-1}\left|v_{\varepsilon, \alpha}\right| d x\right)^{\frac{2^{\sharp}-2}{2 \sharp-1}}\left(\int_{B}\left|v_{\varepsilon, \alpha}\right|^{2^{\sharp}} d x\right)^{\frac{1}{2 \sharp-1}} \\
& =o\left(\varepsilon^{n-4}+\left\|v_{\varepsilon, \alpha}\right\|_{2^{\sharp}}^{2}\right)
\end{aligned}
$$

Now, writing $v_{\varepsilon, \alpha}=-T_{\varepsilon}+\left(v_{\varepsilon, \alpha}+T_{\varepsilon}\right)$ and noting that $v_{\varepsilon, \alpha}+T_{\varepsilon} \in H_{2,0}^{2}(B)$, it comes with Sobolev's inequality that

$$
\begin{equation*}
\left\|v_{\varepsilon, \alpha}\right\|_{2^{\sharp}}=O\left(\left\|\Delta v_{\varepsilon, \alpha}\right\|_{2}+\varepsilon^{\frac{n-4}{2}}\right), \tag{13}
\end{equation*}
$$

and then,

$$
\int_{B} U_{\varepsilon}^{2^{\sharp}-2} v_{\varepsilon, \alpha}^{2} d x=o\left(\left\|\Delta v_{\varepsilon, \alpha}\right\|_{2}^{2}+\varepsilon^{n-4}\right)
$$

This proves (11) and our claim.
Step 2: We now estimate

$$
A_{\varepsilon, \alpha}=\frac{1}{2} \int_{B}\left(\Delta u_{\varepsilon, \alpha}\right)^{2} d x-\frac{1}{2^{\sharp}} \int_{B}\left|u_{\varepsilon, \alpha}\right|^{2^{\sharp}} d x .
$$

Since $u_{\varepsilon, \alpha}=U_{\varepsilon}+v_{\varepsilon, \alpha}+W_{\varepsilon, \alpha}$, we get

$$
\begin{aligned}
& \int_{B}\left(\Delta u_{\varepsilon, \alpha}\right)^{2} d x=\int_{B}\left(\Delta U_{\varepsilon}\right)^{2} d x+\int_{B}\left(\Delta v_{\varepsilon, \alpha}\right)^{2} d x+2 \int_{B} \Delta U_{\varepsilon} \Delta v_{\varepsilon, \alpha} d x \\
& +\int_{B}\left(\Delta W_{\varepsilon, \alpha}\right)^{2} d x+2 \int_{B} \Delta\left(U_{\varepsilon}+v_{\varepsilon, \alpha}\right) \Delta W_{\varepsilon, \alpha} d x
\end{aligned}
$$

Thanks to Green's formula,

$$
\begin{aligned}
& \int_{B} U_{\varepsilon}^{2^{\sharp}-1} v_{\varepsilon, \alpha} d x=\int_{B} \Delta^{2} U_{\varepsilon} v_{\varepsilon, \alpha} d x \\
& =\int_{B} \Delta U_{\varepsilon} \Delta v_{\varepsilon, \alpha} d x+\int_{\partial B}\left(\Delta U_{\varepsilon} \frac{\partial v_{\varepsilon, \alpha}}{\partial n}-v_{\varepsilon, \alpha} \frac{\partial \Delta U_{\varepsilon}}{\partial n}\right) d \sigma \\
& =\int_{B} \Delta U_{\varepsilon} \Delta v_{\varepsilon, \alpha} d x-\int_{\partial B}\left(\Delta U_{\varepsilon} \frac{\partial U_{\varepsilon}}{\partial n}-U_{\varepsilon} \frac{\partial \Delta U_{\varepsilon}}{\partial n}\right) d \sigma
\end{aligned}
$$

with (7). Now with (5), we get that

$$
\begin{equation*}
\int_{B} \Delta U_{\varepsilon} \Delta v_{\varepsilon, \alpha} d x=\int_{B} U_{\varepsilon}^{2^{\sharp}-1} v_{\varepsilon, \alpha} d x+c_{n} \varepsilon^{n-4}+o\left(\varepsilon^{n-4}\right) \tag{14}
\end{equation*}
$$

with $c_{n}=4(n-4) \omega_{n-1} a_{n}^{2(n-4)}$.
The following inequality will be useful throughout the paper. For all $p>1$, for all $\theta \in(0, \min (1, p-1)]$, there exists $C_{p, \theta}>0$ such that

$$
\begin{equation*}
\left.\left||x+y|^{p}-|x|^{p}-p\right| x\right|^{p-2} x y \mid \leq C_{p, \theta}\left(|y|^{p}+|x|^{p-1-\theta}|y|^{1+\theta}\right), \tag{15}
\end{equation*}
$$

for all $x, y \in \mathbb{R}$.
It now follows from inequality (15) that

$$
\begin{aligned}
& \left|\left|u_{\varepsilon, \alpha}\right|^{2^{\sharp}}-U_{\varepsilon}^{2^{\sharp}}-2^{\sharp} U_{\varepsilon}^{2^{\sharp}-1}\left(v_{\varepsilon, \alpha}+W_{\varepsilon, \alpha}\right)\right| \\
& \leq C\left(\left|v_{\varepsilon, \alpha}+W_{\varepsilon, \alpha}\right|^{2^{\sharp}}+U_{\varepsilon}^{2^{\sharp}-2}\left|v_{\varepsilon, \alpha}+W_{\varepsilon, \alpha}\right|^{2}\right)
\end{aligned}
$$

Integrating over $B$ and using (9), (13),(11) and (10), it follows that

$$
\begin{aligned}
& \int_{B}\left|u_{\varepsilon, \alpha}\right|^{2^{\sharp}} d x=\int_{B} U_{\varepsilon}^{2^{\sharp}} d x+2^{\sharp} \int_{B} U_{\varepsilon}^{2^{\sharp}-1} W_{\varepsilon, \alpha} d x+2^{\sharp} \int_{B} U_{\varepsilon}^{2^{\sharp}-1} v_{\varepsilon, \alpha} d x \\
& +o\left(\left\|v_{\varepsilon, \alpha}\right\|_{2^{\sharp}}^{2}\right)+O\left(\left\|W_{\varepsilon, \alpha}\right\|_{2^{\sharp}}^{2}\right)+O\left(\int_{B} U_{\varepsilon}^{2^{\sharp}-2} v_{\varepsilon, \alpha}^{2} d x\right) \\
& =\int_{B} U_{\varepsilon}^{2^{\sharp}} d x+2^{\sharp} \int_{B} U_{\varepsilon}^{2^{\sharp}-1} W_{\varepsilon, \alpha} d x+2^{\sharp} \int_{B} U_{\varepsilon}^{2^{\sharp}-1} v_{\varepsilon, \alpha} d x \\
& +o\left(\varepsilon^{n-4}\right)+o\left(\left\|\Delta v_{\varepsilon, \alpha}\right\|_{2}^{2}\right)
\end{aligned}
$$

This equality combined with equalities (14) and (10) leads to

$$
\begin{aligned}
& A_{\varepsilon, \alpha}=\frac{1}{2} \int_{B}\left(\Delta U_{\varepsilon}\right)^{2} d x-\frac{1}{2^{\sharp}} \int_{B} U_{\varepsilon}^{2^{\sharp}} d x+\frac{1}{2} \int_{B}\left(\Delta v_{\varepsilon, \alpha}\right)^{2} d x-\int_{B} U_{\varepsilon}^{2^{\sharp}-1} W_{\varepsilon, \alpha} d x \\
& +\int_{B} \Delta\left(U_{\varepsilon}+v_{\varepsilon, \alpha}\right) \Delta W_{\varepsilon, \alpha} d x+c_{n} \varepsilon^{n-4}+o\left(\varepsilon^{n-4}\right)+o\left(\left\|\Delta v_{\varepsilon, \alpha}\right\|_{2}^{2}\right)
\end{aligned}
$$

Some straightforward computations give that

$$
\int_{B} U_{\varepsilon}^{2^{\sharp}} d x=\int_{\mathbb{R}^{n}} U^{2^{\sharp}} d x+O\left(\varepsilon^{n}\right) .
$$

Now, noting that

$$
\int_{B}\left(\Delta U_{\varepsilon}\right)^{2} d x=\int_{\mathbb{R}^{n}}\left(\Delta U_{\varepsilon}\right)^{2} d x-\varepsilon^{n-4} \int_{\mathbb{R}^{n}-B}\left(\Delta \varepsilon^{\frac{4-n}{2}} U_{\varepsilon}\right)^{2} d x
$$

it comes with (5) that when $n \geq 5$,

$$
\int_{B}\left(\Delta U_{\varepsilon}\right)^{2} d x=\int_{\mathbb{R}^{n}}(\Delta U)^{2} d x-c_{n} \varepsilon^{n-4}+o\left(\varepsilon^{n-4}\right) .
$$

But since

$$
\int_{\mathbb{R}^{n}}(\Delta U)^{2} d x=\int_{\mathbb{R}^{n}} U^{2^{\sharp}} d x=\frac{1}{K_{0}^{\frac{n}{4}}}
$$

we get that

$$
\begin{align*}
& A_{\varepsilon, \alpha}=\frac{2}{n K_{0}^{\frac{n}{4}}}+\frac{c_{n}}{2} \varepsilon^{n-4}+\frac{1}{2} \int_{B}\left(\Delta v_{\varepsilon, \alpha}\right)^{2} d x-\int_{B} U_{\varepsilon}^{2^{\sharp}-1} W_{\varepsilon, \alpha} d x \\
& +\int_{B} \Delta\left(U_{\varepsilon}+v_{\varepsilon, \alpha}\right) \Delta W_{\varepsilon, \alpha} d x+o\left(\left\|\Delta v_{\varepsilon, \alpha}\right\|_{2}^{2}\right)+o\left(\varepsilon^{n-4}\right) \tag{16}
\end{align*}
$$

Step 3: We now estimate

$$
\alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{u_{\varepsilon, \alpha}}{\alpha}\right) d x .
$$

Since $\rho^{\prime}$ is bounded with (2), there exists $C>0$ such that

$$
\begin{equation*}
|\tilde{\rho}(x+y)-\tilde{\rho}(x)-y \rho(x)| \leq C|y|^{2} \tag{17}
\end{equation*}
$$

for all $x, y \in \mathbb{R}$. Hence,

$$
\begin{align*}
& \alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{u_{\varepsilon, \alpha}}{\alpha}\right) d x  \tag{18}\\
& =\alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{U_{\varepsilon}+v_{\varepsilon, \alpha}}{\alpha}+\frac{W_{\varepsilon, \alpha}}{\alpha}\right) d x \\
& =\alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{U_{\varepsilon}+v_{\varepsilon, \alpha}}{\alpha}\right) d x+\alpha^{2^{\sharp}-1} \int_{B} \rho\left(\frac{U_{\varepsilon}+v_{\varepsilon, \alpha}}{\alpha}\right) W_{\varepsilon, \alpha} d x \\
& +O\left(\alpha^{2^{\sharp}-2} \int_{B} W_{\varepsilon, \alpha}^{2} d x\right) \tag{19}
\end{align*}
$$

Now, using the fact that $\rho^{\prime}$ is bounded, (10) and (13), it comes that

$$
\begin{align*}
& \left|\alpha^{2^{\sharp}-1} \int_{B} \rho\left(\frac{U_{\varepsilon}+v_{\varepsilon, \alpha}}{\alpha}\right) W_{\varepsilon, \alpha} d x-\alpha^{2^{\sharp}-1} \int_{B} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) W_{\varepsilon, \alpha} d x\right| \\
& =O\left(\alpha^{2^{\sharp}-2} \int_{B}\left|v_{\varepsilon, \alpha}\right| W_{\varepsilon, \alpha} d x\right) \\
& =O\left(\alpha^{2^{\sharp}-2}\left\|v_{\varepsilon, \alpha}\right\|_{2^{\sharp}}\left\|W_{\varepsilon, \alpha}\right\|_{2^{\sharp}}\right)=o\left(\left\|\Delta v_{\varepsilon, \alpha}\right\|_{2}^{2}\right)+o\left(\varepsilon^{n-4}\right) . \tag{20}
\end{align*}
$$

But

$$
\begin{align*}
& \int_{B} \Delta\left(U_{\varepsilon}+v_{\varepsilon, \alpha}\right) \Delta W_{\varepsilon, \alpha} d x=\int_{B} \Delta^{2}\left(U_{\varepsilon}+v_{\varepsilon, \alpha}\right) W_{\varepsilon, \alpha} d x \\
& =\int_{B}\left(U_{\varepsilon}^{2^{\sharp}-1}+\alpha^{2^{\sharp}-1} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right)\right) W_{\varepsilon, \alpha} d x . \tag{21}
\end{align*}
$$

Now, putting together (16), (19), (20), (21), it comes that

$$
\begin{align*}
& J_{\alpha}\left(u_{\varepsilon, \alpha}\right)=\frac{2}{n K_{0}^{\frac{n}{4}}}+\frac{c_{n}}{2} \varepsilon^{n-4}+\frac{1}{2}\left\|\Delta v_{\varepsilon, \alpha}\right\|_{2}^{2}-\alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{U_{\varepsilon}+v_{\varepsilon, \alpha}}{\alpha}\right) d x \\
& +o\left(\left\|\Delta v_{\varepsilon, \alpha}\right\|_{2}^{2}\right)+o\left(\varepsilon^{n-4}\right) . \tag{22}
\end{align*}
$$

With (2), inequality (17) can be refined as follows: for any $s \in(0,1)$

$$
|\tilde{\rho}(x+y)-\tilde{\rho}(x)-y \rho(x)| \leq C\left(|x|^{s}+|y|^{s}\right)|y|^{2}
$$

for all $x, y \in \mathbb{R}$, where $C$ depends only on $s$ and $\rho$. We then get that

$$
\left|\tilde{\rho}\left(\frac{U_{\varepsilon}+v_{\varepsilon, \alpha}}{\alpha}\right)-\tilde{\rho}\left(\frac{U_{\varepsilon}}{\alpha}\right)-\rho\left(\frac{U_{\varepsilon}}{\alpha}\right)\left(\frac{v_{\varepsilon, \alpha}}{\alpha}\right)\right| \leq C \frac{U_{\varepsilon}^{s}}{\alpha^{s}} \frac{v_{\varepsilon, \alpha}^{2}}{\alpha^{2}}+\frac{v_{\varepsilon, \alpha}^{2+s}}{\alpha^{2+s}} .
$$

Taking $s>0$ small enough, we get with (13), (9) and Hölder and Sobolev inequalities that

$$
\begin{align*}
& \alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{U_{\varepsilon}+v_{\varepsilon, \alpha}}{\alpha}\right) d x=\alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{U_{\varepsilon}}{\alpha}\right) d x \\
& +\alpha^{2^{\sharp}-1} \int_{B} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) v_{\varepsilon, \alpha} d x+o\left(\left\|\Delta v_{\varepsilon, \alpha}\right\|_{2}^{2}\right)+o\left(\varepsilon^{n-4}\right) . \tag{23}
\end{align*}
$$

Through some integrations by parts and using (7) and (12), we get that

$$
\begin{align*}
\alpha^{2^{\sharp}-1} \int_{B} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right)\left(v_{\varepsilon, \alpha}+T_{\varepsilon}\right) d x & =\int_{B}\left(\Delta v_{\varepsilon, \alpha}\right)^{2} d x+n \frac{\partial v_{\varepsilon, \alpha}}{\partial n}(1) \int_{B} \Delta v_{\varepsilon, \alpha} d x \\
& =\int_{B}\left(\Delta v_{\varepsilon, \alpha}\right)^{2} d x-n \omega_{n-1}\left(\frac{\partial v_{\varepsilon, \alpha}}{\partial n}(1)\right)^{2} \tag{24}
\end{align*}
$$

It now follows from (2) that for all $\nu \in(0,1)$, there exists $C_{\nu}>0$ such that $|\rho(s)| \leq C_{\nu}|s|^{1+\nu}$ for all $s \in \mathbb{R}$. We then get with (5) and (12) that

$$
\begin{equation*}
\alpha^{2^{\sharp}-1} \int_{B} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) T_{\varepsilon} d x=o\left(\varepsilon^{n-4}\right) . \tag{25}
\end{equation*}
$$

Putting inequalities $(22),(23),(24),(25)$ all together and using (5), it comes that

$$
\begin{align*}
& J_{\alpha}\left(u_{\varepsilon, \alpha}\right)=\frac{2}{n K_{0}^{\frac{n}{4}}}-\frac{1}{2}\left\|\Delta v_{\varepsilon, \alpha}\right\|_{2}^{2}-\alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{U_{\varepsilon}}{\alpha}\right) d x+\frac{n^{2}-4 n+2}{4} c_{n} \varepsilon^{n-4} \\
& +o\left(\left\|\Delta v_{\varepsilon, \alpha}\right\|_{2}^{2}\right)+o\left(\varepsilon^{n-4}\right) \tag{26}
\end{align*}
$$

Step 4: We now estimate

$$
\alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{U_{\varepsilon}}{\alpha}\right) d x
$$

that is the third term in the RHS of (26). With (5) and some change of variable, we get that

$$
\alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{U_{\varepsilon}}{\alpha}\right) d x=\omega_{n-1} a_{n}^{n} \varepsilon^{\frac{n}{2}} \alpha^{\frac{n}{n-4}} g(\alpha, \varepsilon),
$$

where

$$
\begin{equation*}
g(\alpha, \varepsilon)=\int_{0}^{\frac{\alpha^{\frac{1}{n-4}}}{a_{n} \sqrt{\varepsilon}}} \tilde{\rho}\left(\left(\alpha^{\frac{2}{n-4}} \varepsilon+r^{2}\right)^{\frac{4-n}{2}}\right) r^{n-1} d r \tag{27}
\end{equation*}
$$

For $n \geq 5$, some standard computations lead to (see for instance [AMS])

$$
\begin{align*}
& \frac{\partial g}{\partial \varepsilon}(\alpha, \varepsilon)=-\frac{\alpha^{\frac{n}{n-4}}}{2 a_{n}^{n} \varepsilon^{\frac{n+2}{2}}} \tilde{\rho}\left(\frac{1}{\alpha}\left(\frac{a_{n}^{2} \varepsilon}{1+a_{n}^{2} \varepsilon^{2}}\right)^{\frac{n-4}{2}}\right)  \tag{28}\\
& -\frac{n-4}{2} \alpha^{\frac{2}{n-4}} \int_{0}^{\frac{\alpha^{\frac{1}{n-4}}}{a_{n} \sqrt{\varepsilon}}}\left(\alpha^{\frac{2}{n-4}} \varepsilon+r^{2}\right)^{\frac{2-n}{2}} \rho\left(\left(\alpha^{\frac{2}{n-4}} \varepsilon+r^{2}\right)^{\frac{4-n}{2}}\right) r^{n-1} d r .
\end{align*}
$$

Similarly,

$$
\begin{align*}
& \frac{\partial^{2} g}{\partial \varepsilon^{2}}(\alpha, \varepsilon)=\frac{n+2}{4 a_{n}^{n}} \alpha^{\frac{n}{n-4}} \varepsilon^{-\frac{n+4}{2}} \tilde{\rho}\left(\frac{1}{\alpha}\left(\frac{a_{n}^{2} \varepsilon}{1+a_{n}^{2} \varepsilon^{2}}\right)^{\frac{n-4}{2}}\right) \\
& +\frac{(n-4) \alpha^{\frac{4}{n-4}}}{4 a_{n}^{4} \varepsilon^{4}} \frac{2 a_{n}^{2} \varepsilon^{2}-1}{\left(1+a_{n}^{2} \varepsilon^{2}\right)^{\frac{n-2}{2}}} \rho\left(\frac{1}{\alpha}\left(\frac{a_{n}^{2} \varepsilon}{1+a_{n}^{2} \varepsilon^{2}}\right)^{\frac{n-4}{2}}\right)  \tag{29}\\
& -\frac{(n-4) \alpha^{\frac{4}{n-4}}}{4\left(1+a_{n}^{2} \varepsilon^{2}\right)^{\frac{n-2}{2}}} \rho\left(\frac{1}{\alpha}\left(\frac{a_{n}^{2} \varepsilon}{1+a_{n}^{2} \varepsilon^{2}}\right)^{\frac{n-4}{2}}\right) \\
& +\frac{(n-2)(n-4) \alpha^{\frac{4}{n-4}}}{4} \int_{0}^{\frac{\alpha^{\frac{1}{n-4}} a_{n} \sqrt{\varepsilon}}{}} \frac{\rho\left(\left(\alpha^{\frac{2}{n-4}} \varepsilon+r^{2}\right)^{\frac{4-n}{2}}\right)}{\left(\alpha^{\frac{2}{n-4}} \varepsilon+r^{2}\right)^{\frac{n-2}{2}}} r^{n-3} d r,
\end{align*}
$$

and

$$
\begin{aligned}
& \frac{\partial^{3} g}{\partial \varepsilon^{3}}(\alpha, \varepsilon)=-\frac{(n+2)(n+4)}{8 a_{n}^{n}} \frac{\alpha^{\frac{n}{n-4}}}{\varepsilon^{\frac{n+6}{2}}} \tilde{\rho}\left(\frac{1}{\alpha}\left(\frac{a_{n}^{2} \varepsilon}{1+a_{n}^{2} \varepsilon^{2}}\right)^{\frac{n-4}{2}}\right) \\
& +\frac{(n-4)(n+10)}{8 a_{n}^{4}} \frac{\alpha^{\frac{4}{n-4}}}{\varepsilon^{5}} \rho\left(\frac{1}{\alpha}\left(\frac{a_{n}^{2} \varepsilon}{1+a_{n}^{2} \varepsilon^{2}}\right)^{\frac{n-4}{2}}\right)\left(1+O\left(\varepsilon^{2}\right)\right) \\
& -\frac{(n-4)^{2}}{8 a_{n}^{8-n}} \varepsilon^{\frac{n-14}{2}} \alpha^{\frac{8-n}{n-4}} \rho^{\prime}\left(\frac{1}{\alpha}\left(\frac{a_{n}^{2} \varepsilon}{1+a_{n}^{2} \varepsilon^{2}}\right)^{\frac{n-4}{2}}\right)\left(1+O\left(\varepsilon^{2}\right)\right) \\
& -\frac{(n-4)^{2}(n-2)}{8} \alpha^{\frac{6}{n-4}} \int_{0}^{\frac{\alpha^{\frac{1}{n-4}} a_{n} \sqrt{\varepsilon}}{}} \frac{\rho\left(\left(\alpha^{\frac{2}{n-4} \varepsilon}+r^{2}\right)^{\frac{4-n}{2}}\right)}{\left(\alpha^{\frac{2}{n-4}} \varepsilon+r^{2}\right)^{\frac{n-2}{2}}} r^{n-5} d r .
\end{aligned}
$$

With (2), we get that
for all $\alpha \in(0,1], \varepsilon>0$. Then, using the integral Taylor identity, (27), (28), (29), and (30), we get that for $n \geq 13$,

$$
\begin{align*}
& \alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{U_{\varepsilon}}{\alpha}\right) d x=  \tag{31}\\
& a_{n}^{n} \varepsilon^{\frac{n}{2}} \alpha^{\frac{n}{n-4}}\left(\int_{\mathbb{R}^{n}} \tilde{\rho}\left(|x|^{4-n}\right) d x-\frac{n-4}{2} \varepsilon \alpha^{\frac{2}{n-4}} \int_{\mathbb{R}^{n}} \rho\left(|x|^{4-n}\right)|x|^{2-n} d x\right. \\
& \left.+\frac{(n-2)(n-4)}{8} \alpha^{\frac{4}{n-4}} \varepsilon^{2} \int_{\mathbb{R}^{n}} \rho\left(|x|^{4-n}\right)|x|^{-n} d x+o\left(\alpha^{\frac{4}{n-4}} \varepsilon^{2}\right)\right) \\
& +O\left(\alpha^{\frac{8}{n-4}} \varepsilon^{n-4}\right)
\end{align*}
$$

Step 5: This step is devoted to the estimation of $\int_{B}\left(\Delta v_{\varepsilon, \alpha}\right)^{2} d x$. It comes from (8) that

$$
\begin{aligned}
& \Delta v_{\varepsilon, \alpha}(r)=n \frac{\partial U_{\varepsilon}}{\partial n}(1) \\
& -n \alpha^{2^{\sharp}-1} \int_{0}^{1} s^{n-1}\left[\int_{s}^{r} t^{1-n}\left\{\int_{0}^{t} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) u^{n-1} d u\right\} d t\right] d s
\end{aligned}
$$

When $n \geq 7$, and since $\Delta v_{\varepsilon, \alpha}$ is radially symmetrical, we have that

$$
\begin{aligned}
& \int_{B}\left(\Delta v_{\varepsilon, \alpha}\right)^{2} d x=n \omega_{n-1}\left(\frac{\partial U_{\varepsilon}}{\partial n}(1)\right)^{2} \\
& +\omega_{n-1} \alpha^{\frac{2(n+4)}{(n-4)}} \int_{0}^{1} r^{n-1}\left[\int_{0}^{r} t^{1-n}\left(\int_{0}^{t} s^{n-1} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) d s\right) d t\right]^{2} d r \\
& -n \omega_{n-1} \alpha^{\frac{2(n+4)}{(n-4)}}\left[\int_{0}^{1} r^{n-1} \int_{0}^{r} t^{1-n}\left(\int_{0}^{t} s^{n-1} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) d s\right) d t d r\right]^{2} \\
& =n \omega_{n-1}\left(\frac{\partial U_{\varepsilon}}{\partial n}(1)\right)^{2} \\
& +\omega_{n-1} \alpha^{\frac{2(n+4)}{(n-4)}} \int_{0}^{1} r^{n-1}\left[\int_{0}^{+\infty} t^{1-n}\left(\int_{0}^{t} s^{n-1} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) d s\right) d t\right. \\
& \left.-\int_{r}^{+\infty} t^{1-n}\left(\int_{0}^{t} s^{n-1} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) d s\right) d t\right]^{2} d r \\
& -n \omega_{n-1} \alpha^{\frac{2(n+4)}{(n-4)}}\left[\int _ { 0 } ^ { 1 } r ^ { n - 1 } \left(\int_{0}^{+\infty} t^{1-n}\left(\int_{0}^{t} s^{n-1} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) d s\right) d t\right.\right. \\
& \left.\left.-\int_{r}^{+\infty} t^{1-n}\left(\int_{0}^{t} s^{n-1} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) d s\right) d t\right) d r\right]^{2}
\end{aligned}
$$

Then, still for $n \geq 7$, we get that

$$
\begin{align*}
& \int_{B}\left(\Delta v_{\varepsilon, \alpha}\right)^{2} d x=n \omega_{n-1}\left(\frac{\partial U_{\varepsilon}}{\partial n}(1)\right)^{2} \\
& +\omega_{n-1} \alpha^{\frac{2(n+4)}{(n-4)}} \int_{0}^{1} r^{n-1}\left[\int_{r}^{+\infty} t^{1-n}\left(\int_{0}^{t} s^{n-1} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) d s\right) d t\right]^{2} d r \\
& -n \omega_{n-1} \alpha^{\frac{2(n+4)}{(n-4)}}\left[\int_{0}^{1} r^{n-1} \int_{r}^{+\infty} t^{1-n}\left(\int_{0}^{t} s^{n-1} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right) d s\right) d t d r\right]^{2} \\
& =A_{1}(\alpha, \varepsilon)+A_{2}(\alpha, \varepsilon)-A_{3}(\alpha, \varepsilon) \tag{32}
\end{align*}
$$

We estimate each of the terms seperately. First, with (5), we have by direct calculation that

$$
\begin{equation*}
A_{1}(\alpha, \varepsilon)=\frac{n(n-4)}{4} c_{n} \varepsilon^{n-4}+o\left(\varepsilon^{n-4}\right) \tag{33}
\end{equation*}
$$

We now deal with $A_{3}(\alpha, \varepsilon)$, that is the third term in the RHS. A change of variable gives

$$
\begin{align*}
& A_{3}(\alpha, \varepsilon)=n \omega_{n-1} a_{n}^{2 n+4} \alpha^{\frac{2(n+4)}{n-4}} \varepsilon^{n+2}  \tag{34}\\
& \times\left[\int_{0}^{\frac{1}{a_{n} \sqrt{\varepsilon}}} r^{n-1}\left(\int_{r}^{+\infty} t^{1-n}\left(\int_{0}^{t} s^{n-1} \rho\left(\frac{\left(\varepsilon+s^{2}\right)^{\frac{4-n}{2}}}{\alpha}\right) d s\right) d t\right) d r\right]^{2}
\end{align*}
$$

Now, when $n \geq 9$, there exists $\nu \in(0,1]$ such that $\nu \in\left(\frac{4}{n-4}, \frac{8}{n-4}\right)$. With (2), it then comes that there exists $C>0$ such that $|\rho(s)| \leq C|s|^{1+\nu}$ for any $s \neq 0$. Some computations then lead to

$$
\left|\int_{0}^{t} s^{n-1} \rho\left(\frac{\left(\varepsilon+s^{2}\right)^{\frac{4-n}{2}}}{\alpha}\right) d s\right| \leq \frac{C}{\alpha^{1+\nu}}\left(\mathbf{1}_{t \leq 1} t^{n}+\mathbf{1}_{t>1}\right),
$$

for all $t \geq 0, \alpha, \varepsilon \in(0,1)$. Plugging this expression in (34), we get that

$$
\begin{equation*}
A_{3}(\alpha, \varepsilon)=O\left(\varepsilon^{n} \alpha^{2\left(\frac{8}{n-4}-\nu\right)}\right)=o\left(\varepsilon^{n-4}\right) \tag{35}
\end{equation*}
$$

We now deal with $A_{2}(\alpha, \varepsilon)$. A change of variable gives

$$
\begin{equation*}
A_{2}(\alpha, \varepsilon)=\omega_{n-1} \alpha^{\frac{2(n+4)}{n-4}} a_{n}^{n+4} \varepsilon^{\frac{n+4}{2}} f(\alpha, \varepsilon) \tag{36}
\end{equation*}
$$

where

$$
f(\alpha, \varepsilon)=\int_{0}^{\frac{1}{a_{n} \sqrt{\varepsilon}}} r^{n-1}\left[\int_{r}^{+\infty} t^{1-n}\left(\int_{0}^{t} s^{n-1} \rho\left(\frac{1}{\alpha}\left(\varepsilon+s^{2}\right)^{\frac{4-n}{2}}\right) d s\right) d t\right]^{2} d r .
$$

With Lebesgue's theorem, and when $n \geq 9$, we get that for any $\alpha>0$,

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} f(\alpha, \varepsilon)=\alpha^{-\left(2^{\sharp}-1\right)} \int_{0}^{+\infty} r^{n-1}\left[\int_{r}^{+\infty} t^{1-n}\left(\int_{0}^{t} s^{n-1} \rho\left(s^{4-n}\right) d s\right) d t\right]^{2} d r \tag{37}
\end{equation*}
$$

Similarly to what was done in Step 4 , we compute $\frac{\partial f}{\partial \varepsilon}(\alpha, \varepsilon)$ and we find:

$$
\frac{\partial f}{\partial \varepsilon}(\alpha, \varepsilon)=O\left(\alpha^{-2} \varepsilon^{\frac{n-14}{2}}\right)+O\left(\alpha^{-\frac{n+2}{n-4}}\right)
$$

when $n \geq 11$. We then get that

$$
\begin{equation*}
f(\alpha, \varepsilon)=f(\alpha, 0)+O\left(\alpha^{-\frac{n+2}{n-4}} \varepsilon\right)+O\left(\alpha^{-2} \varepsilon^{\frac{n-12}{2}}\right) \tag{38}
\end{equation*}
$$

as soon as $n \geq 13$. Now, (32), (33), (35), (36), (37) and (38) give

$$
\begin{align*}
& \left\|\Delta v_{\varepsilon, \alpha}\right\|_{2}^{2}= \\
& \omega_{n-1} \alpha^{2^{\sharp}-1} a_{n}^{n+4} \varepsilon^{\frac{n+4}{2}} \int_{0}^{+\infty} r^{n-1}\left(\int_{r}^{+\infty} t^{1-n}\left(\int_{0}^{t} s^{n-1} \rho\left(s^{4-n}\right) d s\right) d t\right)^{2} d r \\
& +\frac{n(n-4)}{4} c_{n} \varepsilon^{n-4}+o\left(\varepsilon^{n-4}\right)+O\left(\alpha^{\frac{16}{n-4}} \varepsilon^{n-4}\right)+O\left(\alpha^{\frac{n+6}{n-4}} \varepsilon^{\frac{n+6}{2}}\right) \tag{39}
\end{align*}
$$

for $n \geq 13$. Combining (1), (26), (31), (39), and using the expressions of $I_{1}(\rho)$, $I_{2}(\rho), I_{3}(\rho)$, we get that

$$
\begin{align*}
& J_{\alpha}\left(u_{\varepsilon, \alpha}\right)=\frac{2}{n K_{0}^{\frac{n}{4}}}+\frac{(n-4)(n-2)^{2} \omega_{n-1} a_{n}^{2(n-4)}}{2} \varepsilon^{n-4}  \tag{40}\\
& -a_{n}^{n} \omega_{n-1} \varepsilon^{\frac{n}{2}} \alpha^{\frac{n}{n-4}}\left[\frac{I_{1}(\rho)}{n}-\frac{I_{2}(\rho)}{2} \varepsilon \alpha^{\frac{2}{n-4}}+\frac{a_{n}^{4} I_{3}(\rho)}{2(n-4)^{5}} \alpha^{\frac{4}{n-4}} \varepsilon^{2}+o\left(\alpha^{\frac{4}{n-4}} \varepsilon^{2}\right)\right] \\
& +o\left(\varepsilon^{n-4}\right)+O\left(\alpha^{\frac{8}{n-4}} \varepsilon^{n-4}\right)+O\left(\alpha^{\left.\frac{n+6}{n-4} \varepsilon^{\frac{n+6}{2}}\right)}\right.
\end{align*}
$$

as soon as $n \geq 13$. Some similar arguments lead to the following estimates in smaller dimensions:

$$
\begin{align*}
& J_{\alpha}\left(u_{\varepsilon, \alpha}\right)=\frac{2}{n K_{0}^{\frac{n}{4}}}+\frac{(n-4)(n-2)^{2} \omega_{n-1} a_{n}^{2(n-4)}}{2} \varepsilon^{n-4} \\
& -a_{n}^{n} \omega_{n-1} \varepsilon^{\frac{n}{2}} \alpha^{\frac{n}{n-4}}\left[\frac{I_{1}(\rho)}{n}+o(1)\right]+o\left(\varepsilon^{n-4}\right)+O\left(\alpha^{\frac{8}{n-4}} \varepsilon^{n-4}\right) \tag{41}
\end{align*}
$$

for $n \geq 9$. If we assume that $n \geq 11$, we obtain that

$$
\begin{align*}
& J_{\alpha}\left(u_{\varepsilon, \alpha}\right)=\frac{2}{n K_{0}^{\frac{n}{4}}}+\frac{(n-4)(n-2)^{2} \omega_{n-1} a_{n}^{2(n-4)}}{2} \varepsilon^{n-4} \\
& -a_{n}^{n} \omega_{n-1} \varepsilon^{\frac{n}{2}} \alpha^{\frac{n}{n-4}}\left[\frac{I_{1}(\rho)}{n}-\frac{I_{2}(\rho)}{2} \varepsilon \alpha^{\frac{2}{n-4}}+o\left(\alpha^{\frac{2}{n-4}} \varepsilon\right)\right] \\
& +o\left(\varepsilon^{n-4}\right)+O\left(\alpha^{\frac{8}{n-4}} \varepsilon^{n-4}\right) \tag{42}
\end{align*}
$$

## 3. Proof of the theorem - Existence statement

We obtain solutions of problem $\left(E_{r}\right)$ thanks to the Mountain-pass lemma of Ambrosetti and Rabinowitz. We use the following statement of the lemma:
Theorem 3.1 ([AmRa]). Let $F \in C^{1}(V, \mathbb{R})$ where $(V,\|\cdot\|)$ is a Banach space. We assume that:
(i) $F(0)=0$,
(ii) $\exists \lambda, R>0$ such that $F(v) \geq \lambda$ for all $v \in V$ such that $\|v\|=R$,
(iii) $\exists v_{0} \in V$ such that $\lim \sup _{t \rightarrow+\infty} F\left(t v_{0}\right)<0$.

We let $t_{0}>0$ large be such that $\left\|t_{0} v_{0}\right\|>R$ and $F\left(t_{0} v_{0}\right)<0$, and
$\beta=\inf _{\gamma \in \Gamma} \sup F(\gamma(t))$, where $\Gamma=\left\{\gamma:[0,1] \rightarrow V\right.$ s.t. $\left.\gamma(0)=0, \gamma(1)=t_{0} v_{0}\right\}$.
Then there exists a sequence $\left(u_{n}\right)$ in $V$ such that

$$
F\left(u_{n}\right) \rightarrow \beta \quad, \quad F^{\prime}\left(u_{n}\right) \rightarrow 0 \text { strongly in } V^{\prime} .
$$

Moreover, we have that $\lambda \leq \beta \leq \sup _{t \geq 0} F\left(t v_{0}\right)$.

In order to prove the existence of radial solution, we consider the space

$$
V=H_{2,0}^{2}(B) \cap\left\{v \in H_{2,0}^{2}(B) / v \circ \sigma=v, \text { for all } \sigma \in O_{n}(\mathbb{R})\right\}
$$

where $O_{n}(\mathbb{R})$ denotes the group of the isometries of the Euclidean $n$-dimensional space $\mathbb{R}^{n}$. We also consider the functional $F=J_{1}$ (where $J_{1}$ was defined in section 2) defined on $V$. Clearly (i) of the theorem is satisfied. With (2), we get that point (ii) is satisfied. Point (iii) is clearly satisfied for all $v_{0} \in V-\{0\}$. Let $v_{0} \in V-\{0\}$. Then, it follows from theorem 3.1 that there exists a sequence $\left(u_{p}\right) \in H_{2,0}^{2}(B)$ such that

$$
\begin{equation*}
J_{1}\left(u_{p}\right) \rightarrow \beta \quad, \quad J_{1}^{\prime}\left(u_{p}\right) \rightarrow 0 \text { strongly in } V^{\prime} \tag{43}
\end{equation*}
$$

when $p \rightarrow+\infty$. Here $0<\beta \leq \sup _{t \geq 0} J_{1}\left(t v_{0}\right)$.
Step 1: We claim that there exists $u \in V$ such that $u_{p} \rightharpoonup u$ weakly in $H_{2,0}^{2}(B)$ when $p \rightarrow+\infty$. With the additionnal property that

$$
u \neq 0 \text { if } \sup _{t \geq 0} J_{1}\left(t v_{0}\right)<\frac{2}{n K_{0}^{\frac{n}{4}}} .
$$

We prove the claim. It follows from standard arguments that $\left(u_{p}\right)$ is bounded in $H_{2,0}^{2}(B)$. Then there exists $u \in H_{2,0}^{2}(B)$ such that $u_{p} \rightharpoonup u$ weakly in $H_{2,0}^{2}(B)$. Clearly $u \in V$. We now assume that

$$
\sup _{t \geq 0} J_{1}\left(t v_{0}\right)<\frac{2}{n K_{0}^{\frac{n}{4}}} .
$$

We prove that $u \not \equiv 0$ by contradiction. We assume that $u_{p} \rightharpoonup 0$ weakly in $H_{2,0}^{2}(B)$. We can assume that $u_{p} \rightarrow 0$ in $L^{q}(B)$ for all $q \in\left(1,2^{\sharp}\right)$. Then with (2), it comes that

$$
\begin{aligned}
& J_{1}\left(u_{p}\right)=\frac{1}{2} \int_{B}\left(\Delta u_{p}\right)^{2}-\frac{1}{2^{\sharp}} \int_{B}\left|u_{p}\right|^{2^{\sharp}} d x+o(1)=\beta+o(1) \\
& \left\langle J_{1}^{\prime}\left(u_{p}\right), u_{p}\right\rangle=\int_{B}\left(\Delta u_{p}\right)^{2}-\int_{B}\left|u_{p}\right|^{2^{\sharp}} d x+o(1)=o(1)
\end{aligned}
$$

These inequalities combined with the optimal Sobolev inequality (4) then lead to

$$
\left(\frac{n}{2} \beta\right)^{\frac{2}{2 \sharp}} \leq K_{0} \frac{n}{2} \beta .
$$

Since $\beta>0$, we get $\beta \geq \frac{2}{n K_{0}^{\frac{n}{4}}}$. A contradiction. Then $u \not \equiv 0$. The claim is proved.
With (43), we get that for all $\varphi \in C_{c}^{\infty}(B)$ radially symmetrical, we have that

$$
\int_{B} \Delta u \Delta \varphi d x=\int_{B}\left(|u|^{2^{\sharp}-2} u+\rho(u)\right) \varphi d x
$$

It then follows by standard arguments (see for instance [Heb1]) that this equality occurs for all $\varphi \in C_{c}^{\infty}(B)$. And then

$$
\Delta^{2} u=|u|^{2^{\sharp}-2} u+\rho(u)
$$

in the distribution sense. It then follows from arguments due to Van der Vorst $[\mathrm{VdV}]$ and Agmon-Douglis-Nirenberg [ADN] that for any $\nu \in(0,1)$,

$$
u \in C^{4, \nu}(\bar{B})
$$

and that

$$
\Delta^{2} u=|u|^{2^{\sharp}-2} u+\rho(u), \text { in } B, \text { and } u=\frac{\partial u}{\partial n}=0 \text { on } \partial B .
$$

Now, proving the first part of theorem 1.1 on the unit ball reduces to find some suitable functions $v_{0} \in V-\{0\}$ such that

$$
\begin{equation*}
\sup _{t \geq 0} J_{1}\left(t v_{0}\right)<\frac{2}{n K_{0}^{\frac{n}{4}}} . \tag{44}
\end{equation*}
$$

Situations for which this inequality holds can be found in [EsRo]. We consider the test-functions introduced in section 2 . We now let $\varepsilon>0$ and consider $U_{\varepsilon}+v_{\varepsilon}$, where $v_{\varepsilon}=v_{\varepsilon, 1}$ and $U_{\varepsilon}+v_{\varepsilon} \in V$ by construction. By standard variational arguments, there exists $t_{\varepsilon} \in(0,+\infty)$ such that

$$
\sup _{t \geq 0} J_{1}\left(t\left(U_{\varepsilon}+v_{\varepsilon}\right)\right)=J_{1}\left(t_{\varepsilon}\left(U_{\varepsilon}+v_{\varepsilon}\right)\right) .
$$

Step 2: We claim that

$$
\begin{equation*}
t_{\varepsilon}=1+o\left(\varepsilon^{\frac{n-4}{2}}\right)+o\left(\left\|\Delta v_{\varepsilon}\right\|_{2}\right) \tag{45}
\end{equation*}
$$

We prove the claim. It follows from the estimates of section 2 that

$$
\begin{aligned}
& \int_{B}\left(\Delta\left(U_{\varepsilon}+v_{\varepsilon}\right)\right)^{2} d x=\int_{\mathbb{R}^{n}}(\Delta U)^{2} d x+o\left(\varepsilon^{\frac{n-4}{2}}\right)+o\left(\left\|\Delta v_{\varepsilon}\right\|_{2}\right) \\
& \int_{B}\left|U_{\varepsilon}+v_{\varepsilon}\right|^{2^{\sharp}} d x=\int_{\mathbb{R}^{n}} U^{2^{\sharp}} d x+o\left(\varepsilon^{\frac{n-4}{2}}\right)+o\left(\left\|\Delta v_{\varepsilon}\right\|_{2}\right) .
\end{aligned}
$$

Then $J_{1}\left(t_{\varepsilon}\left(U_{\varepsilon}+v_{\varepsilon}\right)\right) \geq J_{1}\left(U_{\varepsilon}+v_{\varepsilon}\right)=\frac{2}{n K_{0}^{\frac{n}{4}}}+o(1)$. It then easily follows that $t_{\varepsilon} \nrightarrow 0$ and $t_{\varepsilon} \nrightarrow+\infty$. Up to a subsequence, $t_{\varepsilon} \rightarrow t_{0} \in(0,+\infty)$. Then,

$$
\begin{aligned}
& 0=\frac{d}{d t} J_{1}\left(t\left(U_{\varepsilon}+v_{\varepsilon}\right)\right)_{\left.\right|_{\varepsilon}} \\
& =t_{\varepsilon} \int_{B}\left(\Delta\left(U_{\varepsilon}+v_{\varepsilon}\right)\right)^{2} d x-t_{\varepsilon}^{2^{\sharp}-1} \int_{B}\left|U_{\varepsilon}+v_{\varepsilon}\right|^{2^{\sharp}} d x+o\left(\varepsilon^{\frac{n-4}{2}}\right)+o\left(\left\|\Delta v_{\varepsilon}\right\|_{2}\right)
\end{aligned}
$$

then, $t_{\varepsilon}-t_{\varepsilon}^{2^{\sharp}-1}=o\left(\varepsilon^{\frac{n-4}{2}}\right)+o\left(\left\|\Delta v_{\varepsilon}\right\|_{2}\right)$. It then follows that $t_{0}=1$, and that (45) holds. This proves the claim.

Step 3: we now prove the first assertion of theorem 1.1. We write

$$
u_{\varepsilon, 1}=t_{\varepsilon}\left(U_{\varepsilon}+v_{\varepsilon}\right)=U_{\varepsilon}+v_{\varepsilon}+\left(t_{\varepsilon}-1\right)\left(U_{\varepsilon}+v_{\varepsilon}\right)
$$

Then (10) is satisfied with $W_{\varepsilon, 1}=\left(t_{\varepsilon}-1\right)\left(U_{\varepsilon}+v_{\varepsilon}\right)$. Taking $\alpha=1$, it follows from (40) that inequality (44) is satisfied with $v_{0}=U_{\varepsilon}+v_{\varepsilon}, \varepsilon>0$ small, provided the hypothesis of the existence statement of the theorem. It follows from Step 1 that there exists a solution to the problem $\left(E_{1}\right)$. The first assertion of the theorem easily follows throughout a rescaling argument.

Remark: with (41) and (42), this result can be extended to the case of an open subset of $\mathbb{R}^{n}$, and to smaller dimensions. Of course, we cannot recover that the solutions are radially symmetrical in the general case. However, the following result holds:

Proposition 3.1. Let $\Omega$ be a smooth subset of $\mathbb{R}^{n}$. We assume that $\left(H_{\rho}\right)$ holds, and that one of the following conditions occurs:
(i) $n \geq 9$ and $I_{1}(\rho)>0$,
(ii) $n \geq 11, I_{1}(\rho)=0$ and $I_{2}(\rho)<0$,
(iii) $n \geq 13, I_{1}(\rho)=I_{2}(\rho)=0$ and $I_{3}(\rho)>0$,
then there exists $u \in C^{4}(\bar{B})$ a nonzero function such that

$$
\Delta^{2} u=|u|^{2^{\sharp}-2} u+\rho(u) \text { in } \Omega, \text { and } u=\frac{\partial u}{\partial n}=0 \text { on } \partial \Omega \text {. }
$$

## 4. Blow-up analysis I

This section and the following are devoted to the proof of the second part of theorem 1.1. We assume that for all $\alpha>0$, there exists $\hat{u}_{\alpha} \in C^{4}(\bar{B}(0, \alpha))$ a smooth positive radially symmetrical function such that

$$
\begin{cases}\Delta^{2} \hat{u}_{\alpha}=\left|\hat{u}_{\alpha}\right|^{2^{\sharp}-2} \hat{u}_{\alpha}+\rho\left(\hat{u}_{\alpha}\right) & \text { in } B(0, \alpha) \\ \hat{u}_{\alpha} \not \equiv 0 & \text { on } \partial B(0, \alpha) . \\ \hat{u}_{\alpha}=\frac{\partial \hat{u}_{\alpha}}{\partial n}=0 & \end{cases}
$$

and

$$
\frac{1}{2} \int_{B(0, \alpha)}\left(\Delta \hat{u}_{\alpha}\right)^{2} d x-\frac{1}{2^{\sharp}} \int_{B(0, \alpha)}\left|\hat{u}_{\alpha}\right|^{2^{\sharp}} d x-\int_{B(0, \alpha)} \tilde{\rho}\left(\hat{u}_{\alpha}\right) d x<\frac{2}{n K_{0}^{\frac{n}{4}}} .
$$

Here and in the sequel, $\rho \in C^{\infty}(\mathbb{R})$ and $\rho$ verifies (2). Up to rescaling, there exists $u_{\alpha} \in C^{4}(\bar{B})$ radially symmetrical such that

$$
\begin{cases}\Delta^{2} u_{\alpha}=\left|u_{\alpha}\right|^{2^{\sharp}-2} u_{\alpha}+\alpha^{2^{\sharp}-1} \rho\left(\frac{u_{\alpha}}{\alpha}\right) & \text { in } B \\ u_{\alpha} \equiv 0 & \text { on } \partial B, \\ u_{\alpha}=\frac{\partial u_{\alpha}}{\partial n}=0 & \\ J_{\alpha}\left(u_{\alpha}\right)<\frac{2}{n K_{0}^{\frac{n}{4}}} & \end{cases}
$$

where

$$
J_{\alpha}(u)=\frac{1}{2} \int_{B}\left(\Delta u_{\alpha}\right)^{2} d x-\frac{1}{2^{\sharp}} \int_{B}\left|u_{\alpha}\right|^{2^{\sharp}} d x-\alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{u_{\alpha}}{\alpha}\right) d x .
$$

Step 1: We claim that $u_{\alpha} \rightharpoonup 0$ weakly in $H_{2,0}^{2}(B)$. We prove the claim. It follows from $\left(I_{\alpha}\right)$ that

$$
\begin{align*}
J_{\alpha}\left(u_{\alpha}\right) & =\frac{2}{n} \int_{B}\left|u_{\alpha}\right|^{2^{\sharp}} d x+\frac{\alpha^{2^{\sharp}-1}}{2} \int_{B} \rho\left(\frac{u_{\alpha}}{\alpha}\right) u_{\alpha} d x-\alpha^{2^{\sharp}} \int_{B} \tilde{\rho}\left(\frac{u_{\alpha}}{\alpha}\right) d x \\
& =\frac{2}{n} \int_{B} u_{\alpha}^{2^{\sharp}} d x+o\left(\left\|u_{\alpha}\right\|_{2^{\sharp}}\right)+o(1)<\frac{2}{n K_{0}^{\frac{n}{4}}}, \tag{46}
\end{align*}
$$

then $\left\|u_{\alpha}\right\|_{2^{\sharp}}=O(1)$, and with $\left(I_{\alpha}\right),\left\|u_{\alpha}\right\|_{H_{2,0}^{2}(B)}=O(1)$. Up to a subsequence, we can assume that it goes weakly to $u \in H_{2,0}^{2}(B)$. Passing through the limit in $\left(I_{\alpha}\right)$, we have that $\Delta^{2} u=|u|^{2^{\sharp}-2} u$ in the weak sense. Considering that $u_{\alpha} \rightharpoonup u$ in $L^{2^{\sharp}}(B)$, we get that

$$
\int_{B}|u|^{2^{\sharp}} d x \leq \liminf _{\alpha \rightarrow 0} \int_{B}\left|u_{\alpha}\right|^{2^{\sharp}} d x \leq \frac{1}{K_{0}^{\frac{n}{4}}} .
$$

We used that $u \in H_{2,0}^{2}(B) \subset H_{2,0}^{2}\left(\mathbb{R}^{n}\right)$. We are now left with proving that $u \equiv 0$. We argue by contradiction, and assume that $u \not \equiv 0$. Then, multiplying by $u \in$ $H_{2,0}^{2}(B)$ and integrating, we get with the Sobolev inequality (4) that

$$
\frac{1}{K_{0}} \leq \frac{\int_{\mathbb{R}^{n}}(\Delta u)^{2} d x}{\left(\int_{\mathbb{R}^{n}}|u|^{2^{\sharp}} d x\right)^{\frac{2}{2 \sharp}}}=\frac{\int_{B}(\Delta u)^{2} d x}{\left(\int_{B}|u|^{2^{\sharp}} d x\right)^{\frac{2}{2 \sharp}}}=\left(\int_{B}|u|^{2^{\sharp}} d x\right)^{1-\frac{2}{2^{\sharp}}} \leq \frac{1}{K_{0}} .
$$

In particular equality holds, and $u \in H_{2,0}^{2}\left(\mathbb{R}^{n}\right)$ is an extremal function for the Euclidean Sobolev inequality. It follows from [EFJ], [Lie], [Lio] that $u$ is smooth and that there exist $\lambda \in \mathbb{R}, C \neq 0$ and $\tilde{x} \in \mathbb{R}^{n}$ such that

$$
u(x)=\frac{C}{\left(\lambda^{2}+|x-\tilde{x}|\right)^{\frac{n-4}{2}}} .
$$

A contradiction, since $u$ is zero outside $B$. Then $u \equiv 0$. The claim is proved.
We go on with the study of the sequence $u_{\alpha}$. Using that $|\rho(r)| \leq C|r|$ for some positive constant $C$ and all $r \in \mathbb{R}$ and the system $\left(I_{\alpha}\right)$, we get that

$$
\begin{aligned}
& \int_{B}\left(\Delta u_{\alpha}\right)^{2} d x=\int_{B}\left|u_{\alpha}\right|^{2^{\sharp}} d x+\alpha^{2^{\sharp}-1} \int_{B} \rho\left(\frac{u_{\alpha}}{\alpha}\right) u_{\alpha} d x \\
& =\int_{B}\left|u_{\alpha}\right|^{2^{\sharp}} d x+O\left(\alpha^{2^{\sharp}-2} \int_{B} u_{\alpha}^{2} d x\right) .
\end{aligned}
$$

Now, the standard Sobolev inequality asserts that

$$
\left(\int_{B}\left|u_{\alpha}\right|^{2^{\sharp}} d x\right)^{\frac{2}{2 \sharp}} \leq K_{0} \int_{B}\left(\Delta u_{\alpha}\right)^{2} d x=K_{0} \int_{B}\left|u_{\alpha}\right|^{2^{\sharp}} d x+o\left(\left\|u_{\alpha}\right\|_{2^{\sharp}}^{2}\right),
$$

and then $\int_{B}\left|u_{\alpha}\right|^{2^{\sharp}} d x \geq \frac{1}{K_{0}^{\frac{n}{4}}}+o(1)$. Then with (46), we get that

$$
\begin{equation*}
\int_{B}\left|u_{\alpha}\right|^{2^{\sharp}} d x=\frac{1}{K_{0}^{\frac{n}{4}}}+o(1) . \tag{47}
\end{equation*}
$$

Now, noting that

$$
\int_{B}\left|u_{\alpha}\right|^{2^{\sharp}} d x \leq\left(\sup _{B} u_{\alpha}\right)^{2^{\sharp}-2} \int_{B} u_{\alpha}^{2} d x
$$

and that $u_{\alpha} \rightarrow 0$ in $L^{2}(B)$, we get that $\sup _{B} u_{\alpha} \rightarrow+\infty$. Following [Rob] and [FHR], we now let $x_{\alpha} \in B$ and $\mu_{\alpha}>0$ such that

$$
u_{\alpha}\left(x_{\alpha}\right)=\mu_{\alpha}^{-\frac{n-4}{2}}=\sup _{B} u_{\alpha} \rightarrow+\infty .
$$

For $x \in \mathbb{R}^{n}$, we now define

$$
\bar{u}_{\alpha}(x)=\mu_{\alpha}^{\frac{n-4}{2}} u_{\alpha}\left(x_{\alpha}+\mu_{\alpha} x\right) \text { if } x \in B_{\alpha}=B\left(-\frac{x_{\alpha}}{\mu_{\alpha}}, \frac{1}{\mu_{\alpha}}\right)
$$

and $\bar{u}_{\alpha}(x)=0$ elsewhere. Clearly, $\bar{u}_{\alpha} \in H_{2,0}^{2}\left(\mathbb{R}^{n}\right)$ satisfies the following system:

$$
\begin{cases}\Delta^{2} \bar{u}_{\alpha}=\left|\bar{u}_{\alpha}\right|^{2^{\sharp}-2} \bar{u}_{\alpha}+\left(\mu_{\alpha}^{\frac{n-4}{2}} \alpha\right)^{2^{\sharp}-1} \rho\left(\frac{\bar{u}_{\alpha}}{\alpha \mu_{\alpha}^{\frac{-4}{2}}}\right) & \text { in } B_{\alpha}  \tag{48}\\ \bar{u}_{\alpha} \not \equiv 0 & \\ \bar{u}_{\alpha}=\frac{\partial \bar{u}_{\alpha}}{\partial n}=0 & \text { on } \partial B_{\alpha},\end{cases}
$$

Step 2: We now claim that

$$
\begin{equation*}
\lim _{\alpha \rightarrow 0} \frac{d\left(x_{\alpha}, \partial B\right)}{\mu_{\alpha}}=+\infty \tag{49}
\end{equation*}
$$

We prove this claim by contradiction. Assume that

$$
\lim _{\alpha \rightarrow 0} \frac{d\left(x_{\alpha}, \partial B\right)}{\mu_{\alpha}}=R \in[0,+\infty)
$$

Since

$$
\int_{\mathbb{R}^{n}}\left(\Delta \bar{u}_{\alpha}\right)^{2} d x=\int_{B_{\alpha}}\left(\Delta \bar{u}_{\alpha}\right)^{2} d x=\int_{B}\left(\Delta u_{\alpha}\right)^{2} d x=O(1),
$$

it comes that $\left\|\bar{u}_{\alpha}\right\|_{H_{2,0}^{2}\left(\mathbb{R}^{n}\right)}$ is bounded. Then, up to a subsequence, $\bar{u}_{\alpha} \rightharpoonup \bar{u} \in$ $H_{2,0}^{2}\left(\mathbb{R}^{n}\right)$. It then follows that

$$
\int_{\mathbb{R}^{n}}|\bar{u}|^{2^{\sharp}} d x \leq \liminf _{\alpha \rightarrow 0} \int_{B_{\alpha}}\left|\bar{u}_{\alpha}\right|^{2^{\sharp}} d x=\liminf _{\alpha \rightarrow 0} \int_{B}\left|u_{\alpha}\right|^{2^{\sharp}} d x=\frac{1}{K_{0}^{\frac{n}{4}}} .
$$

Since $u_{\alpha}$ is radially symmetrical, we can assume that $x_{\alpha}=x_{0}-R_{\alpha} \mu_{\alpha} \vec{n}_{x_{0}}$, where $x_{0} \in \partial B, \vec{n}_{x_{0}}$ is the unit outward vector at $x_{0}$ and $R_{\alpha} \rightarrow R, R_{\alpha}>0$. Clearly, for all $K>0$ and all $\tilde{R}<R$, there exists $\alpha_{0}>0$ such that

$$
\Omega_{K, \tilde{R}}=B(0, K) \cap\left\{x \in \mathbb{R}^{n} /\left(x, \vec{n}_{x_{0}}\right)<\tilde{R}\right\} \subset \subset B_{\alpha}
$$

for all $\alpha \in\left(0, \alpha_{0}\right)$. We denote by $\mathcal{P}_{R}$ the open half-plane

$$
\mathcal{P}_{R}=\left\{x \in \mathbb{R}^{n} /\left(x, \vec{n}_{x_{0}}\right)<R\right\} .
$$

For all $\varphi \in C_{c}^{\infty}\left(\mathcal{P}_{R}\right)$, we define $\hat{\varphi}_{\alpha} \in C_{c}^{\infty}(B)$ such that

$$
\varphi(x)=\mu_{\alpha}^{\frac{n-4}{2}} \hat{\varphi}_{\alpha}\left(x_{\alpha}+\mu_{\alpha} x\right),
$$

for all $x \in \mathbb{R}^{n}$. With $\left(I_{\alpha}\right)$ and a change of variable, we have that

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} \Delta \bar{u}_{\alpha} \Delta \varphi d x=\int_{\mathbb{R}^{n}}\left[\left|\bar{u}_{\alpha}\right|^{2^{\sharp}-2} \bar{u}_{\alpha}+\left(\alpha \mu_{\alpha^{\frac{n-4}{2}}}\right)^{2^{\sharp}-1} \rho\left(\frac{\bar{u}_{\alpha}}{\alpha \mu_{\alpha}^{\frac{n-4}{2}}}\right)\right] \varphi d x . \tag{50}
\end{equation*}
$$

Letting $\alpha$ go to 0 , it comes that

$$
\int_{\mathbb{R}^{n}} \Delta \bar{u} \Delta \varphi d x=\int_{\mathbb{R}^{n}}|\bar{u}|^{2^{\sharp}-2} \bar{u} \varphi d x,
$$

for all $\varphi \in \mathcal{D}\left(\mathcal{P}_{R}\right)$. We now claim that $\bar{u}(x)=0$ almost everywhere on $\mathbb{R}^{n}-\overline{\mathcal{P}}_{R}$. Let $x \in \mathbb{R}^{n}$ such that $\left(x, \vec{n}_{x_{0}}\right)>R$. Then, for $\alpha$ small, $x_{\alpha}+\mu_{\alpha} x \notin B$, and $\bar{u}_{\alpha}(x)=0$. Since $\bar{u}_{\alpha}(x) \rightarrow \bar{u}(x)$ almost everywhere, we get that $\bar{u}(x)=0$ almost everywhere on $\left\{x \in \mathbb{R}^{n} /\left(x, \vec{n}_{x_{0}}\right)>R\right\}$. This claim is proved. It then follows that $\bar{u} \in H_{2,0}^{2}\left(\mathcal{P}_{R}\right)$, and that $\int_{\mathbb{R}^{n}}(\Delta \bar{u})^{2} d x=\int_{\mathbb{R}^{n}}|\bar{u}|^{2^{\sharp}} d x$. With some arguments similar to the ones proceeded in the proof of Step 1, we get that $\bar{u} \equiv 0$. We define
$v_{\alpha}(x)=\bar{u}_{\alpha}\left(x+R_{\alpha} \vec{n}_{x_{0}}\right), x_{0}+\mu_{\alpha} x \in B$. Clearly, there exists a diffeomorphism $\varphi_{\alpha}: B(0, R+2) \rightarrow \mathcal{U}_{\alpha}$, where $\mathcal{U}_{\alpha}$ is an open subset of $\mathbb{R}^{n}$, such that for any $x=\left(x_{1}, \ldots, x_{n}\right) \in B(0, R+2)$,

$$
x_{0}+\mu_{\alpha} \varphi_{\alpha}(x) \in B \Leftrightarrow x_{n}<0 .
$$

We now set $\tilde{v}_{\alpha}=v_{\alpha} \circ \varphi_{\alpha}$. Clearly, there exists a second order operator $L_{\alpha}$ on $B(0, R+2)$ such that

$$
\begin{cases}L_{\alpha}^{2} \tilde{v}_{\alpha}=\left|\tilde{v}_{\alpha}\right|^{2^{\sharp}-2} \tilde{v}_{\alpha}+\left(\alpha \mu_{\alpha}^{\frac{n-4}{2}}\right)^{2^{\sharp}-1} \rho\left(\frac{\tilde{v}_{\alpha}}{\alpha \mu_{\alpha}^{\frac{n-4}{2}}}\right) & \text { in } B(0, R+2) \cap\left\{x_{n}<0\right\} \\ \tilde{v}_{\alpha}=\frac{\partial \tilde{v}_{\alpha}}{\partial n}=0 & \text { on } B(0, R+2) \cap\left\{x_{n}=0\right\}\end{cases}
$$

We can write $L_{\alpha}^{2}$ as follows:

$$
L_{\alpha}^{2}=a_{\alpha}^{i j k l} \partial_{i j k l}+P_{\alpha}\left(\nabla, \nabla^{2}, \nabla^{3}\right)
$$

where $P_{\alpha}$ is a polynomial with continuous and uniformly bounded coefficients, and $a_{i j k l}^{\alpha}$ is also continuous and uniformly bounded with respect to $\alpha$. Moreover, we have that

$$
\frac{1}{2}|X|^{4} \leq a_{i j k l}^{\alpha} X_{i} X_{j} X_{k} X_{l} \leq 2|X|^{4}
$$

for all $X \in \mathbb{R}^{n}$. It then follows from Theorem 15.3 of Agmon-Douglis-Nirenberg [ADN] that for all $p>1$, there exists $C_{p}>0$ such that

$$
\left\|\tilde{v}_{\alpha}\right\|_{H_{4}^{p}\left(B(0, R+1) \cap\left\{x_{n}<0\right\}\right)} \leq C_{p}
$$

Here, we have used that $\left|\tilde{v}_{\alpha}\right| \leq 1$. It then follows that, up to a subsequence, $\tilde{v}_{\alpha}$ converges to a continuous function in $C^{0}(\bar{B}(0, R+1))$. But since $\bar{u}_{\alpha} \rightharpoonup 0$ weakly, it easily comes that $\tilde{v}_{\alpha} \rightarrow 0$ in $C^{0}(\bar{B}(0, R+1))$. A contradiction, since $1=\tilde{v}_{\alpha}\left(-R_{\alpha} \vec{n}_{x_{0}}\right)$. This proves (49) and our claim.

Thanks to (48) and (49), it then follows by standard regularity theory that $\bar{u}_{\alpha}$ is bounded in $C_{l o c}^{4, \beta}\left(\mathbb{R}^{n}\right)$, with $\beta \in(0,1)$. Then, there exists $U_{0} \in C^{4}\left(\mathbb{R}^{n}\right)$ such that

$$
\begin{equation*}
\bar{u}_{\alpha} \rightarrow U_{0} \text { in } C_{l o c}^{4}\left(\mathbb{R}^{n}\right) \tag{51}
\end{equation*}
$$

$U_{0}$ verifies that $\Delta^{2} U_{0}=\left|U_{0}\right|^{2^{\sharp}-2} U_{0},\left|U_{0}(x)\right| \leq U_{0}(0)=1$ for all $x \in \mathbb{R}^{n}$. With some arguments similar to the ones proceeded in Step 1, it comes that $U_{0}$ is an extremal for the Sobolev inequality (4). It follows from [Lin], [HeRo] that

$$
U_{0}(x)=U(x)=\left(\frac{a_{n}^{2}}{a_{n}^{2}+|x|^{2}}\right)^{\frac{n-4}{2}}
$$

for all $x \in \mathbb{R}^{n}$, where $U$ was defined in (3).
Step 3: We now claim that

$$
\begin{equation*}
x_{\alpha}=o\left(\mu_{\alpha}\right) \tag{52}
\end{equation*}
$$

We prove this claim by contradiction. We borrow ideas from Faget [Fag]. We assume that there exists $\eta>0$ such that $\frac{\left|x_{\alpha}\right|}{\mu_{\alpha}} \geq \eta$ up to a subsequence. Let $\vec{n}_{0} \in \mathbb{R}^{n}$ such that $\left|\vec{n}_{0}\right|=1$. Up to a rotation, we can assume that $x_{\alpha}=\left|x_{\alpha}\right| \vec{n}_{0}$. We let $N \in \mathbb{N}^{\star}$ and $\sigma$ an isometry of $\mathbb{R}^{n}$ such that $\sigma^{i}\left(\vec{n}_{0}\right) \neq \vec{n}_{0}$ for $1 \leq i<N$ and $\sigma^{N}\left(\vec{n}_{0}\right)=\vec{n}_{0}$. We let $\delta>0$ such that

$$
\begin{equation*}
\delta<\frac{1}{3} \eta \inf _{\substack{i \neq j \\ 0 \leq i, j<N}}\left|\sigma^{i}\left(\vec{n}_{0}\right)-\sigma^{j}\left(\vec{n}_{0}\right)\right| \tag{53}
\end{equation*}
$$

We now define $B_{\alpha}^{i}=B\left(\sigma^{i}\left(x_{\alpha}\right), \delta \mu_{\alpha}\right)$ for all $i=0, \ldots, N-1$. We claim that $B_{\alpha}^{i} \cap B_{\alpha}^{j}=$ $\emptyset$ for all $i \neq j \in[0, N-1)$. We prove this claim by contradiction. We assume that there exist $k \neq l \in[0, N-1)$ such that $B_{\alpha}^{k} \cap B_{\alpha}^{l} \neq \emptyset$. It then follows that

$$
\left|\sigma^{k}\left(x_{\alpha}\right)-\sigma^{l}\left(x_{\alpha}\right)\right|<2 \delta \mu_{\alpha}
$$

Using that $x_{\alpha}=\left|x_{\alpha}\right| \vec{n}_{0}$, it comes that

$$
\eta \mu_{\alpha} \inf _{\substack{i \neq j \\ 0 \leq i, j<N}}\left|\sigma^{i}\left(\vec{n}_{0}\right)-\sigma^{j}\left(\vec{n}_{0}\right)\right| \leq\left|x_{\alpha}\right| \cdot\left|\sigma^{k}\left(\vec{n}_{0}\right)-\sigma^{l}\left(\vec{n}_{0}\right)\right|<2 \delta \mu_{\alpha},
$$

a contradiction with (53). This claim is proved. Now, using that $u_{\alpha}$ is radially symmetrical, we get that

$$
\begin{array}{r}
\int_{B}\left|u_{\alpha}\right|^{2^{\sharp}} d x \geq \int_{\cup_{i=0}^{N-1} B_{\alpha}^{i}}\left|u_{\alpha}\right|^{2^{\sharp}} d x=\sum_{i=0}^{N-1} \int_{B_{\alpha}^{i}}\left|u_{\alpha}\right|^{2^{\sharp}} d x \\
\geq N \int_{B\left(x_{\alpha}, \delta \mu_{\alpha}\right)}\left|u_{\alpha}\right|^{2^{\sharp}} d x=N \int_{B(0, \delta)}\left|\bar{u}_{\alpha}\right|^{2^{\sharp}} d x .
\end{array}
$$

Now, using (51) and (47), it comes that

$$
N \int_{B(0, \delta)} U^{2^{\sharp}} d x \leq \frac{1}{K_{0}^{\frac{n}{4}}}
$$

for all $N \in \mathbb{N}^{\star}$. A contradiction with (3) and $\delta>0$. It then follows that $\frac{\left|x_{\alpha}\right|}{\mu_{\alpha}} \rightarrow 0$, and the claim is proved.

Step 4: We claim that

$$
\begin{equation*}
\left\|u_{\alpha}-U_{\mu_{\alpha}}\right\|_{H_{2}^{2}(B)} \rightarrow 0 \tag{54}
\end{equation*}
$$

where $U_{\mu_{\alpha}}$ is defined in (5). We prove the claim. We introduce a new rescaled function

$$
\tilde{u}_{\alpha}(x)=\mu_{\alpha^{2}}^{\frac{n-4}{2}} u_{\alpha}\left(\mu_{\alpha} x\right) \quad \text { if } x \in B\left(0, \frac{1}{\mu_{\alpha}}\right)
$$

and $\tilde{u}_{\alpha}(x)=0$ elsewhere. Clearly, $\tilde{u}_{\alpha}$ satisfies the following system:

$$
\begin{cases}\Delta^{2} \tilde{u}_{\alpha}=\left|\tilde{u}_{\alpha}\right|^{2^{\sharp}-2} \tilde{u}_{\alpha}+\left(\mu_{\alpha}^{\frac{n-4}{2}} \alpha\right)^{2^{\sharp}-1} \rho\left(\frac{\tilde{u}_{\alpha}}{\alpha \mu_{\alpha}^{\frac{n-4}{2}}}\right) & \text { in } B\left(0, \frac{1}{\mu_{\alpha}}\right) \\ \tilde{u}_{\alpha} \neq 0 & \text { on } \partial B\left(0, \frac{1}{\mu_{\alpha}}\right) \\ \tilde{u}_{\alpha}=\frac{\partial \tilde{u}_{\alpha}}{\partial n}=0 & \end{cases}
$$

and $\tilde{u}_{\alpha}$ is radially symmetrical. It follows from (51), (52) that

$$
\tilde{u}_{\alpha}(0) \rightarrow 1 \text { and } \tilde{u}_{\alpha} \rightarrow U \text { in } C_{l o c}^{4}\left(\mathbb{R}^{n}\right)
$$

Let $R>0$. It then follows that

$$
\int_{B\left(0, R \mu_{\alpha}\right)}\left|u_{\alpha}\right|^{2^{\sharp}} d x=\int_{B(0, R)}\left|\tilde{u}_{\alpha}\right|^{\left.\right|^{\sharp}} d x=\int_{B(0, R)} U^{2^{\sharp}} d x+o(1) .
$$

Now, by dominated convergence,

$$
\lim _{R \rightarrow+\infty} \int_{B(0, R)} U^{2^{\sharp}} d x=\int_{\mathbb{R}^{n}} U^{2^{\sharp}} d x=\frac{1}{K_{0}^{\frac{n}{4}}} .
$$

With (47), it comes that

$$
\int_{B-B\left(0, R \mu_{\alpha}\right)}\left|u_{\alpha}\right|^{2^{\sharp}} d x=\varepsilon(R)+o(1)
$$

where $\lim _{R \rightarrow+\infty} \varepsilon(R)=0$. Similarly,

$$
\int_{B-B\left(0, R \mu_{\alpha}\right)}\left(\Delta u_{\alpha}\right)^{2} d x=\varepsilon(R)+o(1)
$$

where $\lim _{R \rightarrow+\infty} \varepsilon(R)=0$. Now we get that

$$
\begin{aligned}
& \int_{B}\left(\Delta\left(u_{\alpha}-U_{\mu_{\alpha}}\right)\right)^{2} d x \\
& =\int_{B\left(0, R \mu_{\alpha}\right)}\left(\Delta\left(u_{\alpha}-U_{\mu_{\alpha}}\right)\right)^{2} d x+\int_{B-B\left(0, R \mu_{\alpha}\right)}\left(\Delta\left(u_{\alpha}-U_{\mu_{\alpha}}\right)\right)^{2} d x \\
& =\int_{B(0, R)}\left(\Delta\left(\tilde{u}_{\alpha}-U\right)\right)^{2} d x+O\left(\int_{B-B\left(0, R \mu_{\alpha}\right)}\left(\left(\Delta u_{\alpha}\right)^{2}+\left(\Delta U_{\mu_{\alpha}}\right)^{2}\right) d x\right) \\
& =o(1)+\varepsilon(R)
\end{aligned}
$$

with the strong convergence of $\tilde{u}_{\alpha}$ on compact subsets. Consequently,

$$
\int_{B}\left(\Delta\left(u_{\alpha}-U_{\mu_{\alpha}}\right)\right)^{2} d x \rightarrow 0
$$

Now, clearly, $u_{\alpha}-U_{\mu_{\alpha}} \rightarrow 0$ in $H_{1}^{2}(B)$, the Sobolev space of first order. And then,

$$
\left\|u_{\alpha}-U_{\mu_{\alpha}}\right\|_{H_{2}^{2}(B)} \rightarrow 0
$$

The claim is proved.
Now, for $\alpha, \varepsilon>0$, we consider the function $v_{\varepsilon, \alpha} \in C^{4}(\bar{B})$ defined in (7). Following [AMS], we now consider the minimization problem:

$$
\inf _{\substack{0<\varepsilon \leq 1 \\ 0<a \leq 2}} F_{\varepsilon, a}\left(u_{\alpha}\right),
$$

where

$$
\begin{equation*}
F_{\varepsilon, a}\left(u_{\alpha}\right)=\int_{B}\left[\Delta\left(u_{\alpha}-a\left(U_{\varepsilon}+v_{\varepsilon, \alpha}\right)\right)\right]^{2} d x \tag{55}
\end{equation*}
$$

With (9) and (54), it comes that

$$
\begin{equation*}
\lim _{\alpha \rightarrow 0} \inf _{\substack{0<\varepsilon \leq 1 \\ 0<a \leq 2}} F_{\varepsilon, a}\left(u_{\alpha}\right)=0 \tag{56}
\end{equation*}
$$

We choose $\alpha_{0}>0$ such that

$$
\inf _{\substack{0<\varepsilon \leq 1 \\ 0<a \leq 2}} F_{\varepsilon, a}\left(u_{\alpha}\right)<\frac{1}{3} \int_{\mathbb{R}^{n}}(\Delta U)^{2} d x \text { and } \int_{B}\left(\Delta u_{\alpha}\right)^{2} d x>\frac{1}{2} \int_{\mathbb{R}^{n}}(\Delta U)^{2} d x
$$

for all $\alpha \in\left(0, \alpha_{0}\right)$.
Step 5: We claim that this infimum is attained at $\varepsilon_{\alpha}, a_{\alpha}$, and then that $\varepsilon_{\alpha} \rightarrow 0$ and $a_{\alpha} \rightarrow 1$ when $\alpha \rightarrow 0$. We prove the claim. We fix $\alpha \in\left(0, \alpha_{0}\right)$. We let $a_{p} \in(0,2], \varepsilon_{p} \in(0,1]$ such that

$$
\inf _{\substack{0<\varepsilon \leq 1 \\ 0<a \leq 2}} F_{\varepsilon, a}\left(u_{\alpha}\right)=F_{\varepsilon_{p}, a_{p}}\left(u_{\alpha}\right)+o(1),
$$

where $o(1) \rightarrow 0$ when $p \rightarrow+\infty$. If $a_{p} \rightarrow 0$ when $p \rightarrow+\infty$, then

$$
\inf _{\substack{0<\varepsilon \leq 1 \\ 0<a \leq 2}} F_{\varepsilon, a}\left(u_{\alpha}\right)=\int_{B}\left(\Delta u_{\alpha}\right)^{2} d x
$$

A contradiction with the choice of $\alpha$. Then $a_{p} \rightarrow a_{\alpha} \in(0,2]$ when $p \rightarrow+\infty$. If $\varepsilon_{p} \rightarrow 0$, then for all $\delta \in(0,1)$,

$$
\begin{aligned}
\int_{B-B(0, \delta)}\left(\Delta u_{\alpha}\right)^{2} d x & =\int_{B-B(0, \delta)}\left[\Delta\left(u_{\alpha}-a_{p}\left(U_{\varepsilon_{p}}+v_{\varepsilon_{p}, \alpha}\right)\right)\right]^{2} d x+o(1) \\
& \leq \inf _{\substack{0<\varepsilon \leq 1 \\
0<a \leq 2}} F_{\varepsilon, a}\left(u_{\alpha}\right)+o(1) \\
& <\frac{1}{3} \int_{\mathbb{R}^{n}}(\Delta U)^{2} d x+o(1)
\end{aligned}
$$

Letting $p \rightarrow+\infty$, we get that $\int_{B-B(0, \delta)}\left(\Delta u_{\alpha}\right)^{2} d x \leq \frac{1}{3} \int_{\mathbb{R}^{n}}(\Delta U)^{2} d x$. Passing to the limit $\delta \rightarrow 0$, we get a contradiction with the choice of $\alpha$. Then $\varepsilon_{p} \rightarrow \varepsilon_{\alpha} \in(0,1]$ when $p \rightarrow+\infty$. So the infimum is attained at $\varepsilon_{\alpha}, a_{\alpha}$. Assume that $a_{\alpha} \rightarrow 0$. Then

$$
\begin{aligned}
\int_{B}\left(\Delta u_{\alpha}\right)^{2} d x & =F_{\varepsilon_{\alpha}, a_{\alpha}}\left(u_{\alpha}\right)+o(1) \\
& =\inf _{\substack{0<\varepsilon \leq 1 \\
0<a \leq 2}} F_{\varepsilon, a}\left(u_{\alpha}\right)=o(1)
\end{aligned}
$$

when $\alpha \rightarrow 0$. A contradiction. Then $a_{\alpha} \nrightarrow 0$. Assume that $\varepsilon_{\alpha} \rightarrow \varepsilon_{0}>0$. We have that

$$
\int_{B\left(0, R \mu_{\alpha}\right)}\left[\Delta\left(u_{\alpha}-a_{\alpha}\left(U_{\varepsilon_{\alpha}}+v_{\varepsilon_{\alpha}, \alpha}\right)\right)\right]^{2} d x=o(1)
$$

for all $R>0$. Passing through the limit with (51), we get that $\int_{B(0, R)}(\Delta U)^{2} d x=0$ for all $R>0$. A contradiction. Then $\varepsilon_{\alpha} \rightarrow 0$, and $v_{\varepsilon_{\alpha}, \alpha} \rightarrow 0$ in $H_{2,0}^{2}(B)$ when $\alpha \rightarrow 0$ (see (9)). Now, with (55) and (56),

$$
\begin{aligned}
\int_{B}\left(\Delta u_{\alpha}\right)^{2} d x & =\left(\lim _{\alpha \rightarrow 0} a_{\alpha}\right)^{2} \int_{\mathbb{R}^{n}}(\Delta U)^{2} d x+o(1) \\
& =\frac{\left(\lim _{\alpha \rightarrow 0} a_{\alpha}\right)^{2}}{K_{0}^{\frac{n}{4}}}+o(1)
\end{aligned}
$$

But with $\left(I_{\alpha}\right)$ and (47), we get that $\int_{B}\left(\Delta u_{\alpha}\right)^{2} d x=\int_{B}\left|u_{\alpha}\right|^{2^{\sharp}} d x=\frac{1}{K_{0}^{\frac{n}{4}}}+o(1)$. Consequently, $a_{\alpha} \rightarrow 1$ when $\alpha \rightarrow 0$. This proves our claim.

We now write

$$
\begin{equation*}
u_{\alpha}=a_{\alpha}\left(U_{\varepsilon_{\alpha}}+v_{\varepsilon_{\alpha}, \alpha}\right)+w_{\alpha} . \tag{57}
\end{equation*}
$$

Clearly $w_{\alpha} \rightarrow 0$ in $H_{2,0}^{2}(B)$. For the sake of simplicity, we now write $\varepsilon=\varepsilon_{\alpha} \rightarrow 0$ and $v_{\alpha}=v_{\varepsilon_{\alpha}, \alpha}$. Differentiating $F_{\varepsilon, a}\left(u_{\alpha}\right)$ with respect to $\varepsilon$ and $a$, we get that

$$
\begin{align*}
& \int_{B} \Delta w_{\alpha} \Delta\left(U_{\varepsilon}+v_{\alpha}\right) d x=0,  \tag{58}\\
& \int_{B} \Delta w_{\alpha} \Delta \frac{\partial}{\partial \varepsilon}\left(U_{\varepsilon}+v_{\alpha}\right) d x=0 . \tag{59}
\end{align*}
$$

Next section is devoted to obtaining asymptotic estimates on $\left\|w_{\alpha}\right\|_{H_{2,0}^{2}(B)} \rightarrow 0$ and $1-a_{\alpha} \rightarrow 0$.

## 5. Blow-up analysis II

We follow the techniques developed in [AMS].
Step 1: Some integrations by parts and (58) lead to

$$
\begin{aligned}
& \int_{B}\left(U_{\varepsilon}+v_{\alpha}\right) \Delta^{2} u_{\alpha} d x=\int_{B} \Delta\left(U_{\varepsilon}+v_{\alpha}\right) \Delta u_{\alpha} d x \\
& =a_{\alpha} \int_{B} \Delta\left(U_{\varepsilon}+v_{\alpha}\right) \Delta\left(U_{\varepsilon}+v_{\alpha}\right) d x \\
& =a_{\alpha} \int_{B}\left(U_{\varepsilon}+v_{\alpha}\right)\left(\Delta^{2} U_{\varepsilon}+\Delta^{2} v_{\alpha}\right) d x
\end{aligned}
$$

Using $\left(I_{\alpha}\right),(7)$ and (6) we get that

$$
\begin{align*}
& \int_{B}\left(a U_{\varepsilon}^{2^{\sharp}-1}-\left|u_{\alpha}\right|^{2^{\sharp}-2} u_{\alpha}\right)\left(U_{\varepsilon}+v_{\alpha}\right) d x \\
& =\alpha^{2^{\sharp}-1} \int_{B}\left[\rho\left(\frac{u_{\alpha}}{\alpha}\right)-a_{\alpha} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right)\right]\left(U_{\varepsilon}+v_{\alpha}\right) d x . \tag{60}
\end{align*}
$$

Clearly, for all $p>1$ there exist $C_{p}>0$ such that

$$
\begin{equation*}
\left||x+y|^{p-1}(x+y)-|x|^{p-1} x\right| \leq C_{p}\left(|y|^{p}+|x|^{p-1}|y|\right), \tag{61}
\end{equation*}
$$

for all $x, y \in \mathbb{R}$. Inequalities (11), (61), (13), the definition (57), and some Hölder inequalities lead to

$$
\begin{align*}
& \int_{B}\left(\left|u_{\alpha}\right|^{2^{\sharp}-2} u_{\alpha}-a_{\alpha} U_{\varepsilon}^{2^{\sharp}-1}\right)\left(U_{\varepsilon}+v_{\alpha}\right) d x \\
& =\left(a_{\alpha}^{2^{\sharp}-1}-a_{\alpha}\right) \int_{B} U_{\varepsilon}^{2^{\sharp}-1}\left(U_{\varepsilon}+v_{\alpha}\right) d x \\
& +\int_{B}\left(\left|u_{\alpha}\right|^{2^{\sharp}-2} u_{\alpha}-\left(a_{\alpha} U_{\varepsilon}\right)^{2^{\sharp}-1}\right)\left(U_{\varepsilon}+v_{\alpha}\right) d x \\
& =\left(2^{\sharp}-2\right)\left(a_{\alpha}-1\right) \int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}} d x+O\left(\left\|w_{\alpha}\right\|_{2^{\sharp}}\right)+o\left(\left\|\Delta v_{\alpha}\right\|_{2}\right) \\
& +o\left(\varepsilon^{\frac{n-4}{2}}\right)+o\left(a_{\alpha}-1\right) \tag{62}
\end{align*}
$$

Now, since $\rho^{\prime}$ is bounded, there exists $C>0$ such that $|\rho(x)-\rho(y)| \leq C|x-y|$ for all $x, y \in \mathbb{R}$. It then comes that

$$
\begin{aligned}
& \left|\alpha^{2^{\sharp}-1} \int_{B}\left[\rho\left(\frac{u_{\alpha}}{\alpha}\right)-a_{\alpha} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right)\right]\left(U_{\varepsilon}+v_{\alpha}\right) d x\right| \\
& \leq \alpha^{2^{\sharp}-1} \int_{B}\left|\rho\left(\frac{u_{\alpha}}{\alpha}\right)-\rho\left(\frac{U_{\varepsilon}}{\alpha}\right)\right|\left(U_{\varepsilon}+v_{\alpha}\right) d x \\
& \quad+\left|a_{\alpha}-1\right| \int_{B}\left|\rho\left(\frac{U_{\varepsilon}}{\alpha}\right)\right|\left(U_{\varepsilon}+v_{\alpha}\right) d x \\
& \leq C \alpha^{2^{\sharp}-2} \int_{B}\left|u_{\alpha}-U_{\varepsilon}\right| \times\left|U_{\varepsilon}+v_{\alpha}\right| d x+o\left(a_{\alpha}-1\right) \\
& \leq C \alpha^{2^{\sharp}-2} \int_{B}\left|\left(a_{\alpha}-1\right) U_{\varepsilon}+a v_{\alpha}+w_{\alpha}\right| \times\left|U_{\varepsilon}+v_{\alpha}\right| d x+o\left(a_{\alpha}-1\right) \\
& =o\left(a_{\alpha}-1\right)+o\left(\left\|v_{\alpha}\right\|_{2^{\sharp}}\right)+o\left(\left\|w_{\alpha}\right\|_{2^{\sharp}}\right)
\end{aligned}
$$

This inequality combined with (13), (60) and (62) then gives that

$$
\begin{equation*}
a_{\alpha}-1=O\left(\left\|w_{\alpha}\right\|_{2^{\sharp}}\right)+o\left(\left\|\Delta v_{\alpha}\right\|_{2}\right)+o\left(\varepsilon^{\frac{n-4}{2}}\right) . \tag{63}
\end{equation*}
$$

Step 2: We go on with the estimates of $\left\|w_{\alpha}\right\|_{H_{2,0}^{2}(B)}$ and $1-a_{\alpha}$. First note that for all $p>1$, for all $\theta \in(0, \min \{1, p-1\}]$, there exists $C_{p, \theta}>0$ such that

$$
\left.\left||x+y|^{p-1}(x+y)-|x|^{p-1} x-p\right| x\right|^{p-1} y \mid \leq C_{p, \theta}\left(|y|^{p}+|x|^{p-1-\theta}|y|^{1+\theta}\right),
$$

for any $x, y \in \mathbb{R}$. This inequality and Hölder's inequality give that

$$
\begin{aligned}
& \int_{B} w_{\alpha}\left|u_{\alpha}\right|^{2^{\sharp}-2} u_{\alpha} d x-\int_{B} w_{\alpha}\left(a_{\alpha} U_{\varepsilon}\right)^{2^{\sharp}-1} d x=\left(2^{\sharp}-1\right) \int_{B} U_{\varepsilon}^{2^{\sharp}-2} w_{\alpha}^{2} d x \\
& +o\left(\left\|v_{\alpha}\right\|_{2^{\sharp}}^{2}\right)+o\left(\left\|w_{\alpha}\right\|_{2^{\sharp}}^{2}\right)+\left(2^{\sharp}-1\right) a_{\alpha}^{2^{\sharp}-1} \int_{B} U_{\varepsilon}^{2^{\sharp}-2} v_{\alpha} w_{\alpha} d x .
\end{aligned}
$$

With some more Hölder inequality, we get that

$$
\begin{align*}
& \int_{B} w_{\alpha}\left|u_{\alpha}\right|^{2^{\sharp}-2} u_{\alpha} d x-\int_{B} w_{\alpha}\left(a_{\alpha} U_{\varepsilon}\right)^{2^{\sharp}-1} d x=\left(2^{\sharp}-1\right) \int_{B} U_{\varepsilon}^{2^{\sharp}-2} w_{\alpha}^{2} d x \\
& +o\left(\left\|v_{\alpha}\right\|_{2^{\sharp}}^{2}\right)+o\left(\left\|w_{\alpha}\right\|_{2^{\sharp}}^{2}\right)+O\left(\left\|w_{\alpha}\right\|_{2^{\sharp}}\left\|U_{\varepsilon}\right\|_{2^{\sharp}}^{\frac{2}{}_{\sharp}}-2\right.  \tag{64}\\
& \int_{B} U_{\varepsilon}^{2^{\sharp}-2} v_{\alpha}^{2} d x
\end{align*} .
$$

Some integrations by parts and (58) give that

$$
\begin{aligned}
\int_{B}\left(\Delta w_{\alpha}\right)^{2} d x & =\int_{B} \Delta w_{\alpha} \Delta\left(u_{\alpha}-a_{\alpha}\left(U_{\varepsilon}+v_{\alpha}\right)\right) d x \\
& =\int_{B} \Delta w_{\alpha} \Delta u_{\alpha} d x=\int_{B} w_{\alpha} \Delta^{2} u_{\alpha} d x
\end{aligned}
$$

and that

$$
\int_{B} w_{\alpha} \Delta^{2}\left(U_{\varepsilon}+v_{\alpha}\right) d x=\int_{B} \Delta w_{\alpha} \Delta\left(U_{\varepsilon}+v_{\alpha}\right) d x=0 .
$$

And then

$$
\begin{align*}
\int_{B}\left(\Delta w_{\alpha}\right)^{2} d x= & \int_{B} w_{\alpha} \Delta^{2} u_{\alpha} d x-a_{\alpha}^{2^{\sharp}-1} \int_{B} w_{\alpha} \Delta^{2}\left(U_{\varepsilon}+v_{\alpha}\right) d x \\
= & \int_{B} w_{\alpha}\left[u_{\alpha}^{2^{\sharp}-1} d x-\left(a_{\alpha} U_{\varepsilon}\right)^{2^{\sharp}-1}\right] d x \\
& +\alpha^{2^{\sharp}-1} \int_{B} w_{\alpha}\left[\rho\left(\frac{u_{\alpha}}{\alpha}\right)-a_{\alpha}^{2^{\sharp}-1} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right)\right] d x \tag{65}
\end{align*}
$$

Similarly to what was done in Step 1, we get that

$$
\begin{aligned}
& \alpha^{2^{\sharp}-1} \int_{B} w_{\alpha}\left[\rho\left(\frac{u_{\alpha}}{\alpha}\right)-a_{\alpha}^{2^{\sharp}-1} \rho\left(\frac{U_{\varepsilon}}{\alpha}\right)\right] d x \\
& =o\left(\left|a_{\alpha}-1\right|\left\|w_{\alpha}\right\|_{H_{2,0}^{2}(B)}\right)+o\left(\left\|w_{\alpha}\right\|_{2^{\sharp}}^{2}\right)+o\left(\left\|v_{\alpha}\right\|_{2^{\sharp}}^{2}\right) .
\end{aligned}
$$

Plugging (64) and this latest equality in (65), and using (6) and (11), it comes that

$$
\begin{aligned}
\int_{B}\left(\Delta w_{\alpha}\right)^{2} d x= & \left(2^{\sharp}-1\right) \int_{B} U_{\varepsilon}^{2^{\sharp}-2} w_{\alpha}^{2} d x+o\left(\left|a_{\alpha}-1\right|\left\|w_{\alpha}\right\|_{H_{2,0}^{2}(B)}\right)+o\left(\left\|w_{\alpha}\right\|_{2^{\sharp}}^{2}\right) \\
& +o\left(\left\|\Delta v_{\alpha}\right\|_{2}^{2}\right)+o\left(\varepsilon^{n-4}\right) .
\end{aligned}
$$

Now, with (63), it comes that

$$
\begin{align*}
\int_{B}\left(\Delta w_{\alpha}\right)^{2} d x= & \left(2^{\sharp}-1\right) \int_{B} U_{\varepsilon}^{2^{\sharp}-2} w_{\alpha}^{2} d x+o\left(\left\|w_{\alpha}\right\|_{2^{\sharp}}^{2}\right) \\
& +o\left(\left\|\Delta v_{\alpha}\right\|_{2}^{2}\right)+o\left(\varepsilon^{n-4}\right) . \tag{66}
\end{align*}
$$

We now define $\tilde{w}_{\alpha} \in D_{2}^{2}\left(\mathbb{R}^{n}\right)$ such that $\tilde{w}_{\alpha}(x)=w_{\alpha}(x)$ on $B$ and $\tilde{w}_{\alpha}(x)=0$ elsewhere. We define

$$
C_{\varepsilon}=\frac{\int_{\mathbb{R}^{n}} \Delta \tilde{w}_{\alpha} \Delta U_{\varepsilon} d x}{\left\|\Delta U_{\varepsilon}\right\|_{2}^{2}}, C_{\varepsilon}^{\prime}=\frac{\int_{\mathbb{R}^{n}} \Delta \tilde{w}_{\alpha} \Delta \frac{\partial U_{\varepsilon}}{\partial \varepsilon} d x}{\left\|\Delta \frac{\partial U_{\varepsilon}}{\partial \varepsilon}\right\|_{2}^{2}}
$$

Noting that $\tilde{w}_{\alpha}$ is radially symmetrical, we have that $\int_{\mathbb{R}^{n}} \Delta \tilde{w}_{\alpha} \Delta \partial_{i} U_{\varepsilon} d x=0$ for all $i=1 \ldots n$. It then follows that $\tilde{w}_{\alpha}-C_{\varepsilon} U_{\varepsilon}-C_{\varepsilon}^{\prime} \frac{\partial U_{\varepsilon}}{\partial \varepsilon}$ is orthogonal to to the space spanned by $U_{\varepsilon}, \partial_{\varepsilon} U_{\varepsilon}, \partial_{i} U_{\varepsilon}, i=1 \ldots n$. It then follows from proposition 6.1 of section 6 that

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}\left[\Delta\left(\tilde{w}_{\alpha}-C_{\varepsilon} U_{\varepsilon}-C_{\varepsilon}^{\prime} \partial_{\varepsilon} U_{\varepsilon}\right)\right]^{2} d x \geq \lambda_{3} \int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2}\left(\tilde{w}_{\alpha}-C_{\varepsilon} U_{\varepsilon}-C_{\varepsilon}^{\prime} \partial_{\varepsilon} U_{\varepsilon}\right)^{2} d x \tag{67}
\end{equation*}
$$

where $\lambda_{3}>2^{\sharp}-1$ is independant of $\alpha$. We develop the RHS term, and get

$$
\begin{aligned}
& \int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2}\left(\tilde{w}_{\alpha}-C_{\varepsilon} U_{\varepsilon}-C_{\varepsilon}^{\prime} \partial_{\varepsilon} U_{\varepsilon}\right)^{2} d x=\int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2} \tilde{w}_{\alpha}^{2} d x+C_{\varepsilon}^{2} \int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}} d x \\
&+C_{\varepsilon}^{\prime 2} \int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2}\left(\partial_{\varepsilon} U_{\varepsilon}\right)^{2} d x-2 C_{\varepsilon} \int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}}-1 \tilde{w}_{\alpha} d x-2 C_{\varepsilon}^{\prime} \int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2} \partial_{\varepsilon} U_{\varepsilon} \tilde{w}_{\alpha} d x \\
&+2 C_{\varepsilon} C_{\varepsilon}^{\prime} \int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-1} \partial_{\varepsilon} U_{\varepsilon} d x
\end{aligned}
$$

Clearly, with (6) and (58),

$$
C_{\varepsilon}=\frac{\int_{\mathbb{R}^{n}} \Delta \tilde{w}_{\alpha} \Delta U_{\varepsilon} d x}{\left\|\Delta U_{\varepsilon}\right\|_{2}^{2}}=-\frac{\int_{B} \Delta w_{\alpha} \Delta v_{\alpha} d x}{\left\|\Delta U_{\varepsilon}\right\|_{2}^{2}}=O\left(\left\|v_{\alpha}\right\|_{H_{2}^{2}(B)}\left\|w_{\alpha}\right\|_{H_{2,0}^{2}(B)}\right) .
$$

But as already noticed, $v_{\alpha} \rightarrow 0$ in $H_{2}^{2}(B)$, so $C_{\varepsilon}=o\left(\left\|w_{\alpha}\right\|_{H_{2,0}^{2}(B)}\right)$. Then,

$$
\begin{gathered}
\int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2}\left(\tilde{w}_{\alpha}-C_{\varepsilon} U_{\varepsilon}-C_{\varepsilon}^{\prime} \partial_{\varepsilon} U_{\varepsilon}\right)^{2} d x=\int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2} \tilde{w}_{\alpha}^{2} d x+o\left(\left\|w_{\alpha}\right\|_{H_{2,0}^{2}(B)}^{2}\right) \\
+C_{\varepsilon}^{\prime 2} \int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2}\left(\partial_{\varepsilon} U_{\varepsilon}\right)^{2} d x-2 C_{\varepsilon}^{\prime} \int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2} \partial_{\varepsilon} U_{\varepsilon} \tilde{w}_{\alpha} d x
\end{gathered}
$$

With the equation verified by $\partial_{\varepsilon} U_{\varepsilon}$ (see (71)), its expression in (6), the expression of $v_{\alpha}$ in (7) and (59), we get that

$$
C_{\varepsilon}^{\prime 2} \int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2}\left(\partial_{\varepsilon} U_{\varepsilon}\right)^{2} d x=\frac{\left(\int_{B} \Delta w_{\alpha} \Delta \partial_{\varepsilon} v_{\alpha} d x\right)^{2}}{\left(2^{\sharp}-1\right) \int_{B}\left(\Delta \frac{\partial U_{\varepsilon}}{\partial \varepsilon}\right)^{2} d x}=o\left(\left\|w_{\alpha}\right\|_{H_{2,0}^{2}(B)}^{2}\right) .
$$

Since $\left\|U_{\varepsilon}\right\|_{2^{\sharp}}$ is bounded, we get with Hölder inequality that

$$
\left.\begin{array}{rl}
C_{\varepsilon}^{\prime} \int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2} \partial_{\varepsilon} U_{\varepsilon} \tilde{w}_{\alpha} d x & =O\left(\left\|w_{\alpha}\right\|_{H_{2,0}^{2}(B)} C_{\varepsilon}^{\prime} \sqrt{\int_{\mathbb{R}^{n}} U_{\varepsilon}^{2 \sharp}-2}\left(\partial_{\varepsilon} U_{\varepsilon}\right)^{2} d x\right.
\end{array}\right)
$$

Consequently,

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2}\left(\tilde{w}_{\alpha}-C_{\varepsilon} U_{\varepsilon}-C_{\varepsilon}^{\prime} \partial_{\varepsilon} U_{\varepsilon}\right)^{2} d x=\int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2} \tilde{w}_{\alpha}^{2} d x+o\left(\left\|w_{\alpha}\right\|_{H_{2,0}^{2}(B)}^{2}\right) . \tag{68}
\end{equation*}
$$

Similarly, we get that

$$
\begin{equation*}
\int_{\mathbb{R}^{n}}\left[\Delta\left(\tilde{w}_{\alpha}-C_{\varepsilon} U_{\varepsilon}-C_{\varepsilon}^{\prime} \partial_{\varepsilon} U_{\varepsilon}\right)\right]^{2} d x=\int_{\mathbb{R}^{n}}\left(\Delta \tilde{w}_{\alpha}\right)^{2} d x+o\left(\left\|w_{\alpha}\right\|_{H_{2,0}^{2}(B)}^{2}\right) . \tag{69}
\end{equation*}
$$

Plugging (68) and (69) into (67), we obtain that

$$
\int_{\mathbb{R}^{n}}\left(\Delta \tilde{w}_{\alpha}\right)^{2} d x \geq \lambda_{3} \int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2} \tilde{w}_{\alpha}^{2} d x+o\left(\left\|w_{\alpha}\right\|_{H_{2,0}^{2}(B)}^{2}\right)
$$

with $\lambda_{3}>2^{\sharp}-1$. Now, plugging this inequality into (66), and using (63), we get that

$$
\begin{aligned}
& \left\|w_{\alpha}\right\|_{H_{2,0}^{2}(B)}=o\left(\left\|\Delta v_{\alpha}\right\|_{2}\right)+o\left(\varepsilon^{\frac{n-4}{2}}\right) \\
& 1-a_{\alpha}=o\left(\left\|\Delta v_{\alpha}\right\|_{2}\right)+o\left(\varepsilon^{\frac{n-4}{2}}\right)
\end{aligned}
$$

Step 3: We now prove the last part of theorem 1.1. It then follows from the estimate (40) of section 2 that

$$
\begin{aligned}
& J_{\alpha}\left(u_{\alpha}\right)=\frac{2}{n K_{0}^{\frac{n}{4}}} \\
& -a_{n}^{n} \omega_{n-1} \varepsilon^{\frac{n}{2}} \alpha^{\frac{n}{n-4}}\left[\frac{I_{1}(\rho)}{n}-\frac{I_{2}(\rho)}{2} \varepsilon \alpha^{\frac{2}{n-4}}+\frac{a_{n}^{4} I_{3}(\rho)}{2(n-4)^{5}} \alpha^{\frac{4}{n-4}} \varepsilon^{2}+o\left(\alpha^{\frac{4}{n-4}} \varepsilon^{2}\right)\right] \\
& +\frac{(n-4)(n-2)^{2} \omega_{n-1} a_{n}^{2(n-4)}}{2} \varepsilon^{n-4} \\
& +o\left(\varepsilon^{n-4}\right)+O\left(\alpha^{\frac{8}{n-4}} \varepsilon^{n-4}\right)+O\left(\alpha^{\frac{n+6}{n-4}} \varepsilon^{\frac{n+6}{2}}\right) .
\end{aligned}
$$

when $n \geq 13$. With $\left(I_{\alpha}\right)$, it comes that $J_{\alpha}\left(u_{\alpha}\right) \leq \frac{2}{n K_{0}^{\frac{n}{4}}}$. The last part of the theorem then follows from the study of the three different cases. This completes the proof of theorem 1.1.
Remark: it follows from (41) and (42) that when $n \geq 9$, we get that $I_{1}(\rho) \geq 0$. If $n \geq 11$ and $I_{1}(\rho)=0$, then $I_{2}(\rho) \leq 0$. If $n \geq 13$ and $I_{1}(\rho)=I_{2}(\rho)=0$, then $I_{3}(\rho) \geq 0$.

## 6. A FOURTH ORDER EIGENVALUE PROBLEM ON $\mathbb{R}^{n}$

This section is devoted to the proof of the following proposition:
Proposition 6.1. We consider the following eigenvalue problem:

$$
\Delta^{2} u=\lambda U_{\varepsilon}^{2^{\sharp}-2} u \text { on } H_{2,0}^{2}\left(\mathbb{R}^{n}\right) .
$$

The first eigenvalue is $\lambda=1$, and its eigenspace is the one-dimensional space spanned by $U_{\varepsilon}$. The second eigenvalue is $2^{\sharp}-1$. Its eigenspace is $(n+1)$-dimensional space spanned by $\partial_{\varepsilon} U_{\varepsilon},\left(\partial_{i} U_{\varepsilon}\right)_{i=1, \ldots, n}$. The third eigenvalue is $\lambda_{3}>2^{\sharp}-1$ and is independant of $\varepsilon>0$. More, for all $u \in H_{2,0}^{2}\left(\mathbb{R}^{n}\right)$, the following inequality holds

$$
\int_{\mathbb{R}^{n}}(\Delta u)^{2} d x \geq \lambda_{3} \int_{\mathbb{R}^{n}} U_{\varepsilon}^{2^{\sharp}-2} u^{2} d x
$$

as soon as

$$
\int_{\mathbb{R}^{n}} \Delta u \Delta U_{\varepsilon} d x=\int_{\mathbb{R}^{n}} \Delta u \Delta \partial_{i} U_{\varepsilon} d x=\int_{\mathbb{R}^{n}} \Delta u \Delta \partial_{\varepsilon} U_{\varepsilon} d x=0
$$

for all $i=1, \ldots, n$.
We first consider the function $U_{0}(x)=\left(\frac{1}{1+|x|^{2}}\right)^{\frac{n-4}{2}}$. We let $\lambda \in \mathbb{R}$ and $\varphi \in$ $H_{2,0}^{2}\left(\mathbb{R}^{n}\right)$ such that

$$
\begin{equation*}
\Delta^{2} \varphi=\lambda U_{0}^{2^{\sharp}-2} \varphi . \tag{70}
\end{equation*}
$$

By standard elliptic theory, it comes that $\varphi \in C^{\infty}\left(\mathbb{R}^{n}\right)$. We denote by $\mathbb{S}^{n}$ the unit sphere of $\mathbb{R}^{n+1}$, and we consider the stereographic projection on $\mathbb{S}^{n}$, that is

$$
\begin{array}{ccc}
\pi: \mathbb{S}^{n}-\{N\} & \rightarrow & \mathbb{R}^{n} \\
x & \mapsto & \left(\frac{x_{1}}{1-x_{n+1}}, \ldots, \frac{x_{n}}{1-x_{n+1}}\right)
\end{array}
$$

where $N=(0, \ldots, 0,1)$ is the north pole. We denote by $h$ the round metric on $\mathbb{S}^{n}$. The pull-back of $h$ via $\pi$ gives that

$$
\left(\pi^{-1}\right)^{\star} h=\psi^{\frac{4}{n-4}} \xi
$$

where $\xi$ is the Euclidean metric on $\mathbb{R}^{n}$ and $\psi(x)=\left(\frac{2}{1+|x|^{2}}\right)^{\frac{n-4}{4}}$. On $\left(\mathbb{S}^{n}, h\right)$, the Paneitz-Branson operator is

$$
P_{h}^{n}=\Delta_{h}^{2}+c_{n} \Delta_{h}+d_{n}
$$

where $\Delta_{h}=-\operatorname{div}_{h}(\nabla)$ is the Laplace-Beltrami operator on $\mathbb{S}^{n}$ and

$$
c_{n}=\frac{n^{2}-2 n-4}{2} \quad, \quad d_{n}=\frac{n(n-4)\left(n^{2}-4\right)}{16}
$$

Branson [Bra] showed that this operator enjoys the following nice property: for all $u \in C^{\infty}\left(\mathbb{S}^{n}\right)$, we get that

$$
\left(P_{h}^{n} u\right) \circ \pi^{-1}=\frac{1}{\psi^{2^{\sharp}-1}} \Delta^{2}\left(\psi u \circ \pi^{-1}\right) .
$$

Now, for $\tilde{\varphi} \in C_{c}^{\infty}\left(\mathbb{R}^{n}\right)$, we define $\tilde{u}=\frac{\tilde{\varphi} \circ \pi}{\psi \circ \pi} \in C^{\infty}\left(\mathbb{S}^{n}\right)$. It follows from the preceding conformal law that

$$
\int_{\mathbb{S}^{n}} \tilde{u} P_{h}^{n} \tilde{u} d v_{h}=\int_{\mathbb{R}^{n}}(\Delta \tilde{\varphi})^{2} d x
$$

and

$$
\begin{aligned}
\frac{1}{16} \int_{\mathbb{S}^{n}} \tilde{u}^{2} d v_{h} & =\int_{\mathbb{S}^{n}}\left(\frac{U_{0}}{\psi}\right)^{2^{\sharp}-2}\left(\frac{\tilde{\varphi}}{\psi}\right)^{2} d v_{h} \\
& =\int_{\mathbb{R}^{n}} U_{0}^{2^{\sharp}-2} \tilde{\varphi}^{2} d x
\end{aligned}
$$

where $d v_{h}$ denotes the volume element on the standard sphere $\left(\mathbb{S}^{n}, h\right)$. Since $\varphi \in$ $H_{2,0}^{2}\left(\mathbb{R}^{n}\right)$, we let $\varphi_{p} \in C_{c}^{\infty}\left(\mathbb{R}^{n}\right)$ such that $\varphi_{p} \rightarrow \varphi$ in $H_{2,0}^{2}\left(\mathbb{R}^{n}\right)$. We consider

$$
u_{p}(x)=\frac{\varphi_{p} \circ \pi}{\psi \circ \pi} \in C^{\infty}\left(\mathbb{S}^{n}\right)
$$

It follows from the preceding equalities that $u_{p}$ converges to a function $u \in H_{2}^{2}\left(\mathbb{S}^{n}\right)$, and that $u(x)=\frac{\varphi \circ \pi(x)}{\psi \circ \pi(x)}$ for all $x \in \mathbb{S}^{n}-\{N\}$. Here, $H_{2}^{2}\left(\mathbb{S}^{n}\right)$ is the second order Sobolev space obtained by completion of $C^{\infty}\left(\mathbb{S}^{n}\right)$ for the norm

$$
\|v\|_{H_{2}^{2}\left(\mathbb{S}^{n}\right)}^{2}=\int_{\mathbb{S}^{n}}\left(\Delta_{h} v\right)^{2} d v_{h}+\int_{\mathbb{S}^{n}}|\nabla v|_{h}^{2} d v_{h}+\int_{\mathbb{S}^{n}} v^{2} d v_{h} .
$$

We have that

$$
P_{h}^{n} u=\frac{\lambda}{16} u
$$

on $\mathbb{S}^{n}-\{N\}$ in the distribution sense. Now, following what was done in $[\mathrm{HeRo}]$, we take a cut-off function $\eta_{s}, s>0$ such that $\eta_{s} \equiv 0$ on $B_{h}(N, s), \eta_{s} \equiv 1$ in $\mathbb{S}^{n}-B_{h}(N, 2 s),\left\|\nabla^{k} \eta_{s}\right\|_{\infty} \leq C s^{-k}$ for $k=0,1,2$ and where $C$ is independant of $s$. We choose $t \in C^{\infty}\left(\mathbb{S}^{n}\right)$, and we get that $\eta_{s} t \rightarrow t$ in $H_{2}^{2}\left(\mathbb{S}^{n}\right)$. We omit the details that can be found in [HeRo]. As a consequence, we get that

$$
P_{h}^{n} u=\frac{\lambda}{16} u
$$

in $\mathcal{D}^{\prime}\left(\mathbb{S}^{n}\right)$. It follows from standard elliptic theory that $u \in C^{\infty}\left(\mathbb{S}^{n}\right)$. It follows from [DHL] and [HeRo] that there exists $\mu \in \mathbb{R}$ an element of the spectrum of $\Delta_{h}$ such that $\frac{\lambda}{16}=\mu^{2}+c_{n} \mu+d_{n}$. More, the eigenspace associated to $\frac{\lambda}{16}$ is the eigenspace of $\mu$, considered as an eigenvalue of $\Delta_{h}$. Now for $L$ an operator and $i \in \mathbb{N}^{\star}$, we denote by $\lambda_{i}(L)$ the $i^{t h}$ eigenvalue of $L$ and $E_{i}(L)$ the corresponding eigenspace. As stated in Berger-Gauduchon-Mazet [BGM], we have that

$$
\begin{aligned}
& \lambda_{1}\left(\Delta_{h}\right)=0, \quad \operatorname{dim}\left(E_{1}\left(\Delta_{h}\right)\right)=1 \\
& \lambda_{2}\left(\Delta_{h}\right)=n, \quad \operatorname{dim}\left(E_{2}\left(\Delta_{h}\right)\right)=n+1
\end{aligned}
$$

Now, coming back to our initial question, we obtain that

$$
\begin{array}{ll}
\lambda_{1}\left(P_{h}^{n}\right)=d_{n}, & \operatorname{dim}\left(E_{1}\left(P_{h}^{n}\right)\right)=1 \\
\lambda_{2}\left(P_{h}^{n}\right)=n^{2}+n c_{n}+d_{n}=d_{n}\left(2^{\sharp}-1\right), & \operatorname{dim}\left(E_{2}\left(P_{h}^{n}\right)\right)=n+1 \\
\lambda_{3}\left(P_{h}^{n}\right)>d_{n}\left(2^{\sharp}-1\right) &
\end{array}
$$

We now come back to the initial problem. We let $\lambda \in \mathbb{R}$ and $\varphi \in H_{2,0}^{2}\left(\mathbb{R}^{n}\right)$ such that

$$
\Delta^{2} \varphi=\lambda U_{\varepsilon}^{2^{\sharp}-2} \varphi .
$$

We define $\tilde{\varphi}(x)=\varphi\left(a_{n} \varepsilon x\right)$, and then

$$
\Delta^{2} \tilde{\varphi}=16 d_{n} \lambda U_{0}^{2^{\sharp}-2} \tilde{\varphi} .
$$

Consequently, the three first eigenvalues of (70) $\lambda_{1}, \lambda_{2}, \lambda_{3}$ and their corresponding eigenspaces $E_{1}, E_{2}, E_{3}$ verify

$$
\begin{array}{ll}
\lambda_{1}=1, & \operatorname{dim}\left(E_{1}\right)=1 \\
\lambda_{2}=2^{\sharp}-1, & \operatorname{dim}\left(E_{2}\right)=n+1 \\
\lambda_{3}=\frac{\lambda_{3}\left(P_{h}^{n}\right)}{d_{n}}>2^{\sharp}-1 & \text { is independant of } \varepsilon
\end{array}
$$

Now, as easily checked,

$$
\begin{equation*}
\Delta^{2} \partial_{i} U_{\varepsilon}=\left(2^{\sharp}-1\right) U_{\varepsilon}^{2^{\sharp}-2} \partial_{i} U_{\varepsilon} \text { and } \Delta^{2} \partial_{\varepsilon} U_{\varepsilon}=\left(2^{\sharp}-1\right) U_{\varepsilon}^{2^{\sharp}-2} \partial_{\varepsilon} U_{\varepsilon}, \tag{71}
\end{equation*}
$$

for all $i=1, \ldots, n$, and $\partial_{\varepsilon} U_{\varepsilon}, \partial_{i} U_{\varepsilon}(i=1, \ldots, n)$ are linearly independant. Then the eigenspace of $2^{\sharp}-1$ is spanned by these vectors. The one dimensional eigenspace of $\lambda_{1}$ is clearly spanned by $U_{\varepsilon}$.

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