# CRITICAL FUNCTIONS AND OPTIMAL SOBOLEV INEQUALITIES 

FRÉDÉRIC ROBERT

## 1. Introduction and statement of the results

Let $(M, g)$ be a compact Riemannian manifold of dimension $n \geq 3$. We denote by $H_{1}^{2}(M)$ the standard Sobolev space of functions in $L^{2}(M)$ with one derivative in $L^{2}(M)$. As is well known, it follows from Sobolev's embedding theorem that there exist $A, B \in \mathbb{R}$ such that

$$
\begin{equation*}
\left(\int_{M}|u|^{2^{\star}} d v_{g}\right)^{\frac{2}{2^{\star}}} \leq A \int_{M}|\nabla u|_{g}^{2} d v_{g}+B \int_{M} u^{2} d v_{g} \tag{1}
\end{equation*}
$$

for all $u \in H_{1}^{2}(M)$, where $2^{\star}=\frac{2 n}{n-2}$ is critical. The sharp constant $A$ in (1) is $A_{\text {opt }}(M)=K(n, 2)^{-2}$, where

$$
\frac{1}{K(n, 2)^{2}}=\frac{4}{n(n-2) \omega_{n}^{4 / n}}
$$

and $\omega_{n}$ is the volume of the unit $n$-sphere. In 1995, using intricate developments from blow-up analysis, Hebey and Vaugon [HeVa2, HeVa3] proved that the sharp constant $A_{\text {opt }}(M)$ is attained in (1). It follows that there exists $B_{0}(g)>0$ such that

$$
\left(\int_{M}|u|^{2^{\star}} d v_{g}\right)^{\frac{2}{2^{\star}}} \leq K(n, 2)^{2}\left(\int_{M}|\nabla u|_{g}^{2} d v_{g}+B_{0}(g) \int_{M} u^{2} d v_{g}\right) \quad\left(I_{o p t, g}\right)
$$

for all $u \in H_{1}^{2}(M)$. Following Hebey [Heb2], we choose $B_{0}(g)$ such that it cannot be lowered. Natural questions with respect to $\left(I_{o p t, g}\right)$, see Druet and Hebey [DrHe] for a presentation in book form, are to compute $B_{0}(g)$, and to know wether or not ( $I_{\text {opt }, g}$ ) possesses extremal functions. Important notions with respect to these questions are the notions of weakly critical and critical functions introduced by Hebey and Vaugon [HeVa4]. These notions have independant applications and were used by Druet [Dru] to prove the Brézis-Nirenberg conjecture on low dimensions. For $f$ a smooth function on $M$, we let

$$
\mu_{g, f}=\inf \frac{\int_{M}\left(|\nabla u|_{g}^{2}+f u^{2}\right) d v_{g}}{\left(\int_{M}|u|^{2^{\star}} d v_{g}\right)^{\frac{2}{2 \star}}}
$$

where the infimum is taken over the nonzero functions of $H_{1}^{2}(M)$, and $d v_{g}$ is the Riemannian volume element. It is by now standard that $\mu_{g, f} \leq \frac{1}{K(n, 2)^{2}}$. By definition, following Hebey and Vaugon [HeVa4], we say that $f$ is a weakly critical function for $g$ if $\mu_{g, f}=\frac{1}{K(n, 2)^{2}}$, and we say that $f$ is a critical function if it is
weakly critical and such that $\mu_{g, \tilde{f}}<\frac{1}{K(n, 2)^{2}}$ for all smooth functions $\tilde{f}$ satisfying that $\tilde{f} \leq f$ and $\tilde{f} \not \equiv f$. Let

$$
B_{0}(g)_{e x t}=\frac{n-2}{4(n-1)} \max _{M} S_{g}
$$

where $S_{g}$ denotes the scalar curvature of $g$. It is easily checked that when $n \geq 4$, we always have that $B_{0}(g) \geq B_{0}(g)_{\text {ext }}$. Moreover, as proved by Djadli and Druet [ DjDr ], one of the two following conditions hold:
(i) either $B_{0}(g)=B_{0}(g)_{e x t}$,
(ii) or ( $I_{\text {opt }, g}$ ) possesses extremal functions.

We know since Hebey [Heb1] that there are examples of manifolds for which (i) is true and (ii) is false. The examples in [Heb1] are in the conformal class of the unit $n$-sphere, and we are left with the question (question (Q1) in [HeVa4]) of describing the class of manifolds for which (i) is true and (ii) is false. Such a question is closely related to the following conjecture of Hebey and Vaugon [HeVa4]: if $n \geq 6$, the Weyl curvature tensor Weyl ${ }_{g}$ of $g$ vanishes up to the order 1 at some point $x_{0}$, and $f$ is a smooth nonegative function vanishing at $x_{0}$, then there exists $g_{f} \in[g]$ a conformal metric to $g$ such that the scalar curvature $S_{g_{f}}$ is maximal at $x_{0}$ and such that $\alpha=B_{0}\left(g_{f}\right)_{\text {ext }}-f$ is a weakly critical function for $g_{f}$. We prove in this article that this conjecture is true. We also treat the case of the dimension $n=5$. Concerning terminology, we say that the Weyl tensor Weyl $l_{g}$ of $g$ vanishes up to order $p \geq 1$ at $x_{0}$ if $W_{\text {eyl }}^{g}\left(x_{0}\right)=0$ and $\nabla^{j} W_{e y l}\left(x_{0}\right)=0$ for all $1 \leq j \leq p$. Our main result states as follows:
Theorem 1.1. Let $(M, g)$ be a smooth compact Riemannian $n$-manifold, $n \geq 5, x_{0}$ a point in $M$, and $f$ a smooth nonnegative nonzero function on $M$ with the property that $f\left(x_{0}\right)=0$. We assume that Weylg is null up to the order 2 at $x_{0}$ if $n=5$, and up to the order 1 at $x_{0}$ if $n \geq 6$. Then there exists $g_{f} \in[g]$ a conformal metric to $g$ such that $S_{g_{f}}$ is maximal at $x_{0}$ and such that $\alpha=\frac{n-2}{4(n-1)} \max _{M} S_{g_{f}}-f$ is a weakly critical function for $g_{f}$. In particular, for any smooth compact Riemannian manifold $(M, g)$ of dimension $n \geq 5$, whose Weyl curvature tensor satisfies the above conditions, there exists a conformal metric $\tilde{g}$ to $g$ such that $B_{0}(\tilde{g})=B_{0}(\tilde{g})_{\text {ext }}$ and $\left(I_{o p t, \tilde{g}}\right)$ does not possess extremal functions.

Hebey and Vaugon [HeVa4] also conjectured that the following surprising complementary result to Theorem 1.1 should hold: if $(M, g)$ is a smooth compact Riemannian manifold of dimension $n \geq 6$ such that its Weyl tensor is null up to the order 2 at some point, then there exists $\tilde{g} \in[g]$ a conformal metric to $g$ such that $B_{0}(\tilde{g})=B_{0}(\tilde{g})_{\text {ext }}=\operatorname{Vol}_{\tilde{g}}^{-2 / n}$ and such that $\left(I_{o p t, \tilde{g}}\right)$ possesses constant extremal functions. In this statement, $\operatorname{Vol}_{\tilde{g}}$ denotes the volume of $M$ wrt the metric $\tilde{g}$. As a remark, it easily follows from the approach we use to prove Theorem 1.1 that this conjecture is also true. With respect to the mathematics developed in Hebey and Vaugon [HeVa4], both conjectures reduce to the proof that for a blowing-up sequence $\left(u_{\epsilon}\right)$ of solutions of the equations attached to the above problems, if $x_{0}$ is the geometric concentration point of the sequence $\left(u_{\epsilon}\right)$, then $d_{g}\left(x_{0}, x_{\epsilon}\right)=O\left(\mu_{\epsilon}\right)$ where the $x_{\epsilon}$ 's are the centers and the $\mu_{\epsilon}$ 's are the weights of the bubbles developed by the $u_{\epsilon}$ 's.

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## 2. Proof of theorem 1.1

When possible, we follow the approach developed in $[\mathrm{HeVa} 4]$. Let $(M, g)$ be a compact Riemannian manifold of dimension $n \geq 3$. We first assume that $n \geq 6$ and that there exists $x_{0} \in M$ such that $W e y l_{g}\left(x_{0}\right)=0$ and $\nabla W e y l_{g}\left(x_{0}\right)=0$. Up to a conformal change of metric (see for instance Lee-Parker [LePa] and [HeVa1]), we can assume that

$$
\begin{equation*}
R m_{g}\left(x_{0}\right)=0 \text { and } \nabla R m_{g}\left(x_{0}\right)=0 . \tag{2}
\end{equation*}
$$

Here $R m_{g}$ and $\nabla R m_{g}$ denote the Riemann tensor and its covariant derivative. Note that in this case, $\nabla^{i} S_{g}\left(x_{0}\right)=0$ for $i=0 \ldots 2$. We denote by $[g]$ the conformal class of the metric $g$ and

$$
[g]_{s}=\left\{\tilde{g} \in[g] \text { s.t. } S_{\tilde{g}}\left(x_{0}\right)=\max _{M} S_{\tilde{g}}\right\}
$$

We let $f \in C^{\infty}(M)-\{0\}$ such that for any $x \in M$,

$$
\begin{equation*}
f(x) \geq f\left(x_{0}\right)=0 \tag{3}
\end{equation*}
$$

In this section, we prove that there exists $\tilde{g} \in[g]_{s}$ such that $\alpha_{\tilde{g}}=c_{n} \max _{M} S_{\tilde{g}}-f$ is a weakly critical function. Here and in the sequel, we let $c_{n}=\frac{n-2}{4(n-1)}$. We argue by contradiction, and assume that for any $\tilde{g} \in[g]_{s}$,

$$
\inf _{u \in H_{1}^{2}(M) \backslash\{0\}} \frac{\int_{M}\left(|\nabla u|_{\tilde{g}}^{2}+\alpha_{\tilde{g}} u^{2}\right) d v_{\tilde{g}}}{\left(\int_{M}|u|^{2^{\star}} d v_{\tilde{g}}\right)^{\frac{2}{2^{\star}}}}<\frac{1}{K(n, 2)^{2}}
$$

Letting $\varphi \in C^{\infty}(M)$ positive such that $\tilde{g}=\varphi^{\frac{4}{n-2}} g$, the preceding inequality is equivalent to

$$
\begin{equation*}
\inf _{u \in H_{1}^{2}(M) \backslash\{0\}} \frac{\int_{M}\left(|\nabla u|_{g}^{2}+h_{\varphi} u^{2}\right) d v_{g}}{\left(\int_{M}|u|^{2^{\star}} d v_{g}\right)^{\frac{2}{2^{\star}}}}<\frac{1}{K(n, 2)^{2}}, \tag{4}
\end{equation*}
$$

where

$$
\begin{equation*}
h_{\varphi}=c_{n} S_{g}+\varphi^{2^{\star}-2}\left(c_{n}\left(\max _{M} S_{\tilde{g}}-S_{\tilde{g}}\right)-f\right) \tag{5}
\end{equation*}
$$

for all positive functions $\varphi$. Let $\delta \in\left(0, i_{g}(M) / 2\right)$, where $i_{g}(M)$ is the injectivity radius of $(M, g)$, and $\eta \in C^{\infty}(M)$ such that $\eta \equiv 1$ in $B_{x_{0}}(\delta)$ and $\eta \equiv 0$ in $M \backslash B_{x_{0}}(2 \delta)$. We also let $c>0$ and consider the function

$$
\varphi_{\epsilon}(x)=\frac{\eta(x)}{\left(\epsilon^{2}+d_{g}\left(x, x_{0}\right)^{2}\right)^{\frac{n-2}{8}}}+(1-\eta(x)) c .
$$

Clearly, $\varphi_{\epsilon}$ is smooth and positive.
Step 1: It follows from [HeVa4] that the metric $g_{\epsilon}=\varphi_{\epsilon}^{\frac{4}{n-2}} g$ belongs to $[g]_{s}$ for small $\epsilon$. Moreover, noting $h_{\epsilon}=h_{\varphi_{\epsilon}}$, we get that

$$
\begin{equation*}
h_{\epsilon}(x) \geq a(\epsilon) d_{g}\left(x, x_{0}\right)^{2} \tag{6}
\end{equation*}
$$

for all $x \in M$ and all $\epsilon>0$, with

$$
\begin{equation*}
\lim _{\epsilon \rightarrow 0} a(\epsilon)=+\infty \tag{7}
\end{equation*}
$$

It is by now classical that it follows from (4) that there exists $u_{\epsilon} \in C^{\infty}(M)$ positive such that

$$
\Delta_{g} u_{\epsilon}+h_{\epsilon} u_{\epsilon}=\lambda_{\epsilon} u_{\epsilon}^{2^{\star}-1} \text { and } \int_{M} u_{\epsilon}^{2^{\star}} d v_{g}=1
$$

where

$$
\begin{equation*}
\lambda_{\epsilon}=\mu_{g, h_{\epsilon}}<\frac{1}{K(n, 2)^{2}} \tag{8}
\end{equation*}
$$

It then follows that $u_{\epsilon} \rightharpoonup 0$ weakly in $H_{1}^{2}(M)$ when $\epsilon \rightarrow 0$. We refer to [HeVa4] for the details. We let $x_{\epsilon} \in M$ and $\mu_{\epsilon}>0$ such that

$$
\max _{M} u_{\epsilon}=u_{\epsilon}\left(x_{\epsilon}\right)=\mu_{\epsilon}^{-\frac{n-2}{2}}
$$

Since $u_{\epsilon} \rightharpoonup 0$ weakly in $H_{1}^{2}(M)$, it follows from $\left(E_{\epsilon}\right)$ that $\lim _{\epsilon \rightarrow 0} \mu_{\epsilon}=0$. We let $y_{\epsilon} \in M$ such that

$$
\frac{d_{g}\left(x_{\epsilon}, y_{\epsilon}\right)}{\mu_{\epsilon}}=O(1)
$$

when $\epsilon \rightarrow 0$. With (6), it follows from [DrRo], see also [DrHe], that there exists $C>0$ such that

$$
\begin{equation*}
d_{g}\left(y_{\epsilon}, x\right)^{\frac{n-2}{2}} u_{\epsilon}(x) \leq C \text { and } u_{\epsilon}(x) \leq C\left(\frac{\mu_{\epsilon}}{\mu_{\epsilon}^{2}+d_{g}\left(x, y_{\epsilon}\right)^{2}}\right)^{\frac{n-2}{2}} \tag{9}
\end{equation*}
$$

for all $x \in M$ and $\epsilon>0$. Moreover, it follows from [HeVa4] that $y_{\epsilon} \rightarrow x_{0}$ when $\epsilon \rightarrow 0$.

Step 2: We denote by $\exp _{y_{\epsilon}}$ the exponential map at $y_{\epsilon}$ with respect to the metric $g$. We let $\tilde{u}_{\epsilon}(x)=u_{\epsilon}\left(\exp _{y_{\epsilon}}(x)\right)$ for all $x \in B_{0}(3 \delta)$ and $g_{\epsilon}=\exp _{y_{\epsilon}}^{\star} g$ the pullback metric of $g$ via the exponential chart at $y_{\epsilon}$. In the sequel, we denote by $g_{\epsilon}^{i j}$ the coordinates of $g^{-1}$, the inverse of the metric tensor, via the chart $\exp _{y_{\epsilon}}$. We let $\delta \in\left(0, i_{g}(M) / 3\right)$. We let $\eta \in C^{\infty}\left(\mathbb{R}^{n}\right)$ such that $\eta \equiv 1$ in $B_{0}(\delta)$ and $\eta \equiv 0$ in $\mathbb{R}^{n} \backslash B_{0}(2 \delta)$. Th optimal Euclidean Sobolev inequality asserts that

$$
\left(\int_{\mathbb{R}^{n}}|u|^{2^{\star}} d x\right)^{\frac{2}{2^{\star}}} \leq K(n, 2)^{2} \int_{\mathbb{R}^{n}}|\nabla u|_{\xi}^{2} d x
$$

for all smooth compactly supported function $u$ on $\mathbb{R}^{n}$. It follows from this inequality that

$$
\begin{equation*}
\left(\int_{\mathbb{R}^{n}}\left(\eta \tilde{u}_{\epsilon}\right)^{2^{\star}} d x\right)^{\frac{2}{2^{\star}}} \leq K(n, 2)^{2} \int_{\mathbb{R}^{n}}\left|\nabla\left(\eta \tilde{u}_{\epsilon}\right)\right|_{\xi}^{2} d x \tag{10}
\end{equation*}
$$

We denote the volume element by $d v_{g_{\epsilon}}=\sqrt{\left|g_{\epsilon}\right|} d x$. We then get that

$$
\begin{align*}
& \int_{\mathbb{R}^{n}}\left(\eta \tilde{u}_{\epsilon}\right)^{2^{\star}} d x  \tag{11}\\
& =1+\int_{B_{0}(\delta)} \tilde{u}_{\epsilon}^{2^{\star}}\left(1-\sqrt{\left|g_{\epsilon}\right|}\right) d x+\int_{B_{0}(2 \delta) \backslash B_{0}(\delta)}\left(\eta \tilde{u}_{\epsilon}\right)^{2^{\star}} d x-\int_{M \backslash B_{y_{\epsilon}}(\delta)} u_{\epsilon}^{2^{\star}} d v_{g}
\end{align*}
$$

It follows from Cartan's expansion of the metric that

$$
g_{\epsilon}^{i j}(x)=\delta^{i j}+O\left(\left|R m_{g}\left(y_{\epsilon}\right)\right|_{g}|x|^{2}+\left|\nabla R m_{g}\left(y_{\epsilon}\right)\right|_{g}|x|^{3}+|x|^{4}\right)
$$

uniformly for $x \in B_{0}(3 \delta)$ when $\epsilon \rightarrow 0$. Plugging this expansion in (11), using ( $E_{\epsilon}$ ) and (9), we get that

$$
\begin{align*}
& \left|\left(\int_{\mathbb{R}^{n}}\left(\eta \tilde{u}_{\epsilon}\right)^{2^{\star}} d v_{\xi}\right)^{\frac{2}{2^{\star}}}-1\right| \leq C\left|R m_{g}\left(y_{\epsilon}\right)\right|_{g} \int_{B_{0}(\delta)} \tilde{u}_{\epsilon}^{2} d x  \tag{12}\\
& +C\left|\nabla R m_{g}\left(y_{\epsilon}\right)\right|_{g} \int_{B_{0}(\delta)}|x| \tilde{u}_{\epsilon}^{2} d x+C \int_{B_{0}(\delta)}|x|^{2} \tilde{u}_{\epsilon}^{2} d x+C \int_{M \backslash B_{y_{\epsilon}}(\delta)} \tilde{u}_{\epsilon}^{2} d x .
\end{align*}
$$

We now deal with the LHS of (10). Clearly,

$$
\begin{aligned}
& \int_{\mathbb{R}^{n}}\left|\nabla\left(\eta \tilde{u}_{\epsilon}\right)\right|_{\xi}^{2} d v_{\xi}=\int_{\mathbb{R}^{n}} g_{\epsilon}^{i j} \partial_{i}\left(\eta \tilde{u}_{\epsilon}\right) \partial_{j}\left(\eta \tilde{u}_{\epsilon}\right) \sqrt{\left|g_{\epsilon}\right|} d x \\
& +\int_{\mathbb{R}^{n}} g_{\epsilon}^{i j} \partial_{i}\left(\eta \tilde{u}_{\epsilon}\right) \partial_{j}\left(\eta \tilde{u}_{\epsilon}\right)\left(1-\sqrt{\left|g_{\epsilon}\right|}\right) d x+\int_{\mathbb{R}^{n}}\left(\delta^{i j}-g_{\epsilon}^{i j}\right) \partial_{i}\left(\eta \tilde{u}_{\epsilon}\right) \partial_{j}\left(\eta \tilde{u}_{\epsilon}\right) d x
\end{aligned}
$$

Noting $\eta_{\epsilon}=\eta \circ \exp _{y_{\epsilon}}^{-1}$, integrating the first term of the RHS by parts and using Cartan's expansion of the metric, we then get that

$$
\begin{aligned}
& \int_{\mathbb{R}^{n}}\left|\nabla\left(\eta \tilde{u}_{\epsilon}\right)\right|_{\xi}^{2} d v_{\xi}=\int_{M} \eta_{\epsilon}^{2} u_{\epsilon} \Delta_{g} u_{\epsilon} d v_{g}+\int_{M}\left|\nabla \eta_{\epsilon}\right|_{g}^{2} u_{\epsilon}^{2} d v_{g} \\
& +O\left(\int_{M}\left(\left|R m_{g}\left(y_{\epsilon}\right)\right|_{g} r^{2}+\left|\nabla R m_{g}\left(y_{\epsilon}\right)\right|_{g} r^{3}+r^{4}\right)\left|\nabla\left(\eta_{\epsilon} u_{\epsilon}\right)\right|_{g}^{2} d v_{g}\right)
\end{aligned}
$$

where $r=d_{g}\left(x, y_{\epsilon}\right)$. Integrating by parts the last term of the RHS, using equation $\left(E_{\epsilon}\right)$, the estimate (9) and that $h_{\epsilon} \geq 0$, we get that

$$
\begin{align*}
& \int_{\mathbb{R}^{n}}\left|\nabla\left(\eta \tilde{u}_{\epsilon}\right)\right|_{\xi}^{2} d v_{\xi} \leq \lambda_{\epsilon}-\lambda_{\epsilon} \int_{M} \eta_{\epsilon}^{2} h_{\epsilon} u_{\epsilon}^{2} d v_{g}  \tag{13}\\
& +C\left|R m_{g}\left(y_{\epsilon}\right)\right|_{g} \int_{B_{y_{\epsilon}(\delta)}} u_{\epsilon}^{2} d v_{g}+C\left|\nabla R m_{g}\left(y_{\epsilon}\right)\right|_{g} \int_{B_{y_{\epsilon}(\delta)}} r u_{\epsilon}^{2} d v_{g} \\
& +C \int_{B_{y_{\epsilon}(\delta)}(\delta} r^{2} u_{\epsilon}^{2} d v_{g}+C \int_{M \backslash B_{y_{\epsilon}}(\delta)} u_{\epsilon}^{2} d v_{g},
\end{align*}
$$

in this expression, $r=d_{g}\left(x, y_{\epsilon}\right)$. Plugging (12) and (13) into (10) and using (8), we get that

$$
\begin{aligned}
& \int_{B_{y_{\epsilon}}(\delta)} h_{\epsilon} u_{\epsilon}^{2} d v_{g} \leq C\left|R m_{g}\left(y_{\epsilon}\right)\right|_{g} \int_{B_{y_{\epsilon}}(\delta)} u_{\epsilon}^{2} d v_{g}+C\left|\nabla R m_{g}\left(y_{\epsilon}\right)\right|_{g} \int_{B_{y_{\epsilon}( }(\delta)} r u_{\epsilon}^{2} d v_{g} \\
& +C \int_{B_{y_{\epsilon}}(\delta)} r^{2} u_{\epsilon}^{2} d v_{g}+C \int_{M \backslash B_{y_{\epsilon}}(\delta)} u_{\epsilon}^{2} d v_{g}
\end{aligned}
$$

in this expression, $r=d_{g}\left(x, y_{\epsilon}\right)$. It follows from [HeVa4] that

$$
\int_{M \backslash B_{y_{\epsilon}}(\delta)} u_{\epsilon}^{2} d v_{g}=o\left(\int_{B_{y_{\epsilon}}(\delta)} h_{\epsilon} u_{\epsilon}^{2} d v_{g}\right) .
$$

It then follows from these last two estimates that

$$
\begin{aligned}
& \int_{B_{y_{\epsilon}}(\delta)} h_{\epsilon} u_{\epsilon}^{2} d v_{g} \leq C\left|R m_{g}\left(y_{\epsilon}\right)\right|_{g} \int_{B_{y_{\epsilon}}(\delta)} u_{\epsilon}^{2} d v_{g} \\
& +C\left|\nabla R m_{g}\left(y_{\epsilon}\right)\right|_{g} \int_{B_{y_{\epsilon}}(\delta)} d_{g}\left(x, y_{\epsilon}\right) u_{\epsilon}^{2} d v_{g}+C \int_{B_{y_{\epsilon}}(\delta)} d_{g}\left(x, y_{\epsilon}\right)^{2} u_{\epsilon}^{2} d v_{g}
\end{aligned}
$$

Since $d_{g}\left(x, y_{\epsilon}\right)^{2} \leq 2 d_{g}\left(x, x_{0}\right)^{2}+2 d_{g}\left(x_{0}, y_{\epsilon}\right)^{2}$, we get with (6) and (7) that

$$
\begin{align*}
\int_{B_{y_{\epsilon}}(\delta)} h_{\epsilon} u_{\epsilon}^{2} d v_{g} \leq & C \cdot\left(d_{g}\left(x_{0}, y_{\epsilon}\right)^{2}+\left|R m_{g}\left(y_{\epsilon}\right)\right|_{g}\right) \cdot \int_{B_{y_{\epsilon}}(\delta)} u_{\epsilon}^{2} d v_{g} \\
& +C\left|\nabla R m_{g}\left(y_{\epsilon}\right)\right|_{g} \int_{B_{y_{\epsilon}}(\delta)} d_{g}\left(x, y_{\epsilon}\right) u_{\epsilon}^{2} d v_{g} \tag{14}
\end{align*}
$$

With (2), we then get that

$$
\begin{align*}
\int_{B_{y_{\epsilon}}(\delta)} h_{\epsilon} u_{\epsilon}^{2} d v_{g} \leq & C d_{g}\left(x_{0}, y_{\epsilon}\right)^{2} \int_{B_{y_{\epsilon}}(\delta)} u_{\epsilon}^{2} d v_{g} \\
& +C d_{g}\left(x_{0}, y_{\epsilon}\right) \int_{B_{y_{\epsilon}( }(\delta)} d_{g}\left(x, y_{\epsilon}\right) u_{\epsilon}^{2} d v_{g} \tag{15}
\end{align*}
$$

as soon as $\frac{d_{g}\left(x_{\epsilon}, y_{\epsilon}\right)}{\mu_{\epsilon}}=O(1)$ when $\epsilon \rightarrow 0$.
Step 3: We claim that

$$
\begin{equation*}
\frac{d_{g}\left(x_{\epsilon}, x_{0}\right)}{\mu_{\epsilon}}=O(1) \tag{16}
\end{equation*}
$$

when $\epsilon \rightarrow 0$. We prove the claim by contradiction and assume that

$$
\begin{equation*}
\lim _{\epsilon \rightarrow 0} \frac{d_{g}\left(x_{\epsilon}, x_{0}\right)}{\mu_{\epsilon}}=+\infty \tag{17}
\end{equation*}
$$

Inequality (15) is obviously verified with $y_{\epsilon}=x_{\epsilon}$. Since $B_{x_{\epsilon}}\left(\mu_{\epsilon}\right) \subset B_{x_{\epsilon}}(\delta)$ when $\epsilon \rightarrow 0$, we get with a change of variable that

$$
\begin{aligned}
& \mu_{\epsilon}^{2} \int_{B_{0}(1)} h_{\epsilon}\left(\exp _{x_{\epsilon}}\left(\mu_{\epsilon} x\right)\right) v_{\epsilon}^{2} d v_{\tilde{g}_{\epsilon}} \leq C d_{g}\left(x_{0}, x_{\epsilon}\right)^{2} \mu_{\epsilon}^{2} \int_{B_{0}\left(\delta \mu_{\epsilon}^{-1}\right)} v_{\epsilon}^{2} d v_{\tilde{g}_{\epsilon}} \\
& +C d_{g}\left(x_{0}, x_{\epsilon}\right) \mu_{\epsilon}^{3} \int_{B_{0}\left(\delta \mu_{\epsilon}^{-1}\right)}|x| v_{\epsilon}^{2} d v_{\tilde{g}_{\epsilon}}
\end{aligned}
$$

where $v_{\epsilon}(x)=\mu_{\epsilon}^{\frac{n-2}{2}} u_{\epsilon}\left(\exp _{y_{\epsilon}}\left(\mu_{\epsilon} x\right)\right)$ for all $x \in B_{0}\left(\delta / \mu_{\epsilon}\right)$ and $\tilde{g}_{\epsilon}=\exp _{x_{\epsilon}}^{\star} g\left(\mu_{\epsilon} x\right)=$ $g_{\epsilon}\left(\mu_{\epsilon} x\right)$. For any $x \in B_{0}(1)$, we get with (17) that

$$
d_{g}\left(x_{0}, \exp _{x_{\epsilon}}\left(\mu_{\epsilon} x\right)\right) \geq d_{g}\left(x_{0}, x_{\epsilon}\right)-\mu_{\epsilon} \geq \frac{1}{2} d_{g}\left(x_{0}, x_{\epsilon}\right)
$$

With the lower bound (6) of $h_{\epsilon}$, we then get that

$$
\begin{aligned}
& a(\epsilon) d_{g}\left(x_{0}, x_{\epsilon}\right)^{2} \mu_{\epsilon}^{2} \int_{B_{0}(1)} v_{\epsilon}^{2} d v_{\tilde{g}_{\epsilon}} \leq C d_{g}\left(x_{0}, x_{\epsilon}\right)^{2} \mu_{\epsilon}^{2} \int_{B_{0}\left(\delta \mu_{\epsilon}^{-1}\right)} v_{\epsilon}^{2} d v_{\tilde{g}_{\epsilon}} \\
& +C d_{g}\left(x_{0}, x_{\epsilon}\right) \mu_{\epsilon}^{3} \int_{B_{0}\left(\delta \mu_{\epsilon}^{-1}\right)}|x| v_{\epsilon}^{2} d v_{\tilde{g}_{\epsilon}} .
\end{aligned}
$$

It follows from Moser's iterative scheme (see for instance [DrRo]) that

$$
\int_{B_{0}(1)} v_{\epsilon}^{2} d v_{\tilde{g}_{\epsilon}} \geq C v_{\epsilon}(0)=C>0
$$

Using the estimates (9) and assuming that $n \geq 6$, we get that

$$
a(\epsilon) d_{g}\left(x_{0}, x_{\epsilon}\right)^{2} \mu_{\epsilon}^{2} \leq C d_{g}\left(x_{0}, y_{\epsilon}\right)^{2} \mu_{\epsilon}^{2}+C d_{g}\left(x_{0}, y_{\epsilon}\right) \mu_{\epsilon}^{3}
$$

A contradiction with (7) and (17). Then (16) holds. This proves the claim.
It follows from (16) that (15) holds with $y_{\epsilon}=x_{0}$. A contradiction, since $h_{\epsilon} \geq$
0 . We have then contradicted our initial assumption (4). Then there exists $\tilde{g}$
conformal to $g$ such that $c_{n} \max _{M} S_{\tilde{g}}-f$ is a weakly critical function. Now we claim that $B_{0}(\tilde{g})$ has no extremal function. We prove the claim by contradiction and assume that $B_{0}(\tilde{g})$ has an extremal function. By definition, $B_{0}(\tilde{g})$ is weakly critical: since it has an extremal function, it is a critical function. Since $B_{0}(\tilde{g}) \geq$ $c_{n} \max _{M} S_{\tilde{g}} \geq c_{n} \max _{M} S_{\tilde{g}}-f$ (see for instance [HeVa4]) and $c_{n} \max _{M} S_{\tilde{g}}-f$ is weakly critical, we get that $B_{0}(\tilde{g})=c_{n} \max _{M} S_{\tilde{g}}=c_{n} \max _{M} S_{\tilde{g}}-f$ and then $f \equiv 0$. A contradiction with the choice of $f$ in (3). Then $B_{0}(\tilde{g})$ has no extremal function. It then follows from $[\mathrm{DjDr}]$ that $B_{0}(\tilde{g})=B_{0}(\tilde{g})_{\text {ext }}$. This proves Theorem 1.1 when $n \geq 6$. Concerning the case of dimension $n=5$, the proof is similar and uses inequality (14). Note also that when $\nabla^{i} W e y l_{g}\left(x_{0}\right)=0$ for $i=0 \ldots 2$, it follows from [HeVa1] that we can assume that $\nabla^{i} R m_{g}\left(x_{0}\right)=0$ for $i=0 \ldots 2$ up to a conformal change of metric. This ends the proof of Theorem 1.1.

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Department of mathematics, ETH Zentrum, CH-8092 Zürich, Switzerland
E-mail address: frederic.robert@math.ethz.ch

